

Lecture #14 of 18(?)

Time-Dependence in Electrochemistry

Chapters 4 and 5

Q: What's in this set of lectures?

A: B&F Chapters 4 & 5 main concepts:

- Section 4.4.2: Fick's Second Law of Diffusion
- Section 5.1: Overview of step experiments
- Section 5.2: Potential step under diffusion controlled
- Sections 5.3 & 5.9: Ultramicroelectrodes
- Sections 5.7 – 5.8: Chronoamperometry/Chronocoulometry

Looking forward... Section 4.4.2 and Chapter 5

- Fick's Second Law of Diffusion
- Linear Diffusion = time-dependent current (Cottrell Equation)
- Anson Plots for surface adsorbed species
- Radial Diffusion = time-independent current (at steady-state)
- **Ultramicroelectrodes (UMEs)**
- **Scanning Electrochemical Microscopy (SECM)**
- **Single molecule electrochemistry**

RECALL:

... so in a potential step experiment...

1. current changes continuously with time.
2. radial diffusion (AKA “edge effects”) limits the data acquisition time window to $\sim 1 - 10$ s.
3. charging imposes a lower limit of $0.1 - 0.5 \mu\text{s}$ on this data acquisition time window.
4. theoretically, maximum current densities are $> 60 \text{ mA cm}^{-2}$ initially, but just $100 \mu\text{A cm}^{-2}$ at $S/N \approx 10$.

... **but, why do we care?**

Why do we care? One reason...

RECALL:

we need to push this up in order to perform meaningful measurements of the kinetics of rapid redox reactions

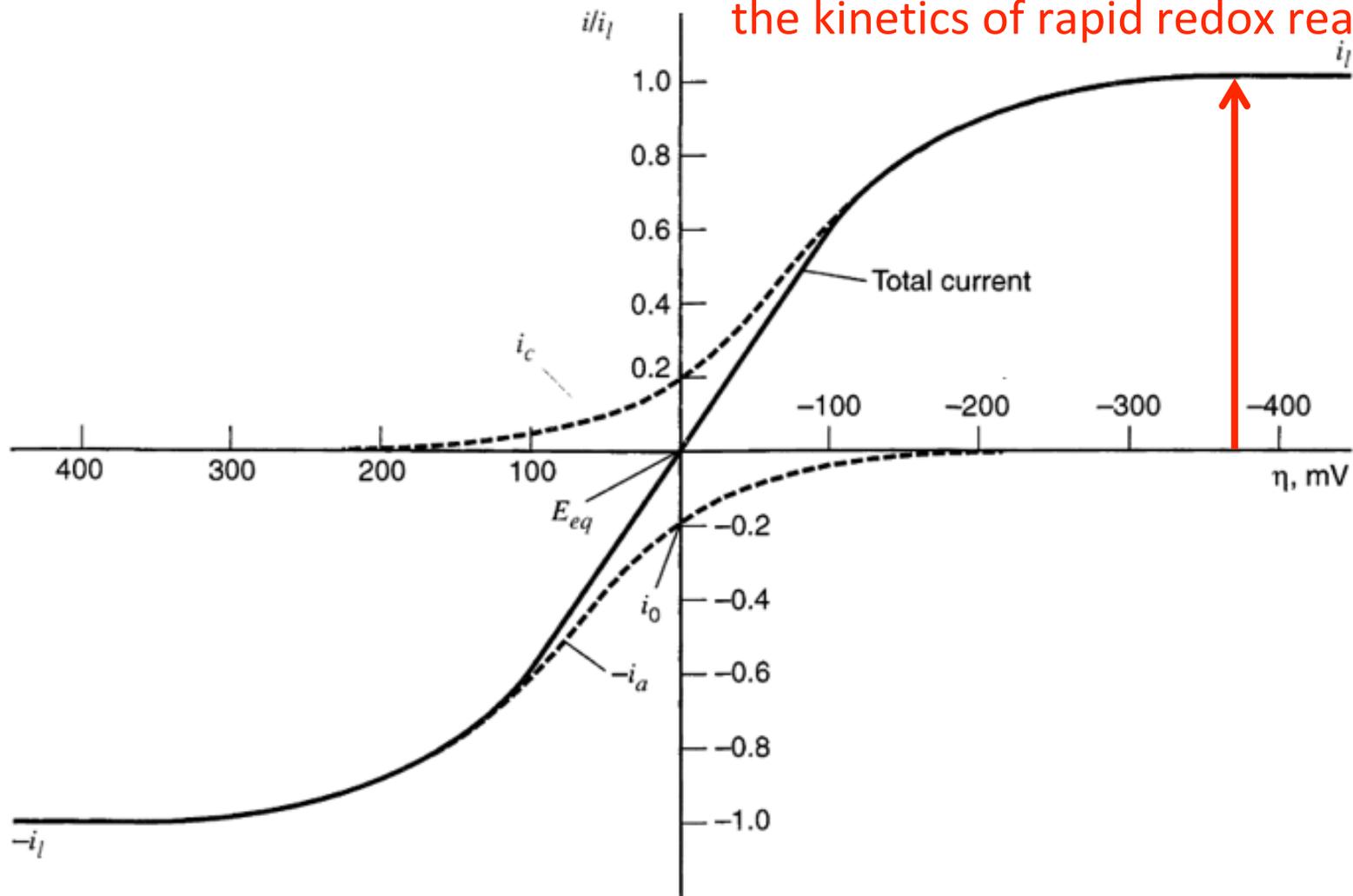


Figure 3.4.1 Current-overpotential curves for the system $O + e \rightleftharpoons R$ with $\alpha = 0.5$, $T = 298$ K, $i_{l,c} = -i_{l,a} = i_l$ and $i_0/i_l = 0.2$. The dashed lines show the component currents i_c and i_a .

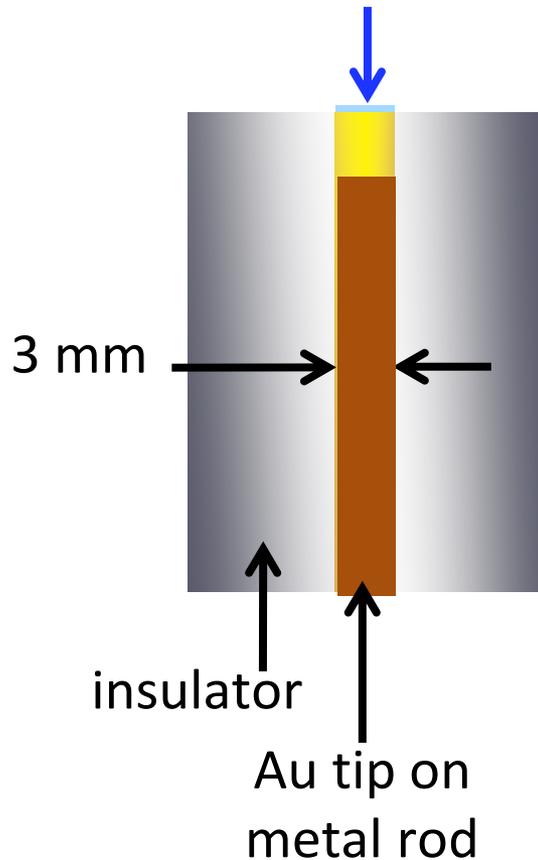
... but deleterious edge effects also suggest an opportunity:

RECALL:

What if instead of avoiding radial diffusion, we exploit it?... Wait, what?

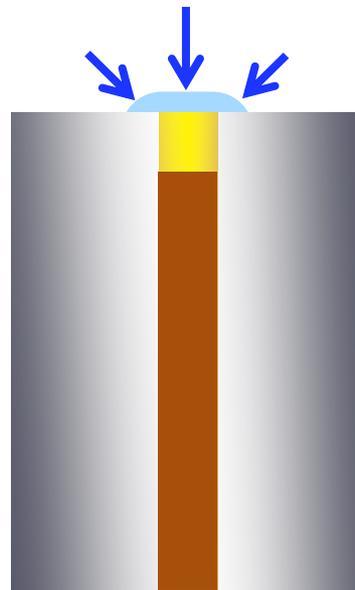
$t \approx 0.1 - 1 \text{ s}$

linear diffusion

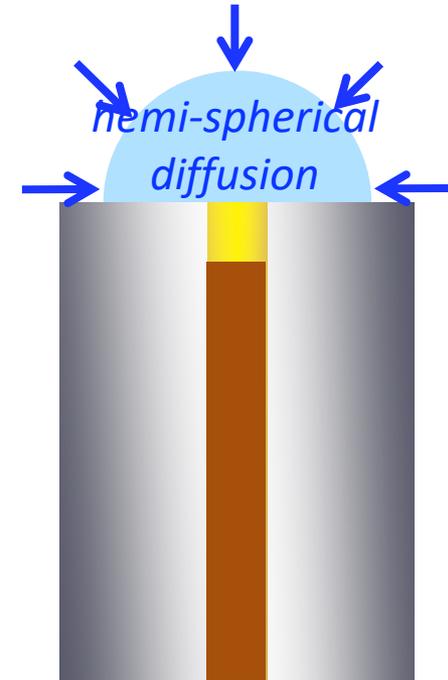


$t \approx 1 - 100 \text{ s}$

mixed diffusion



$t > \sim 100 \text{ s}$

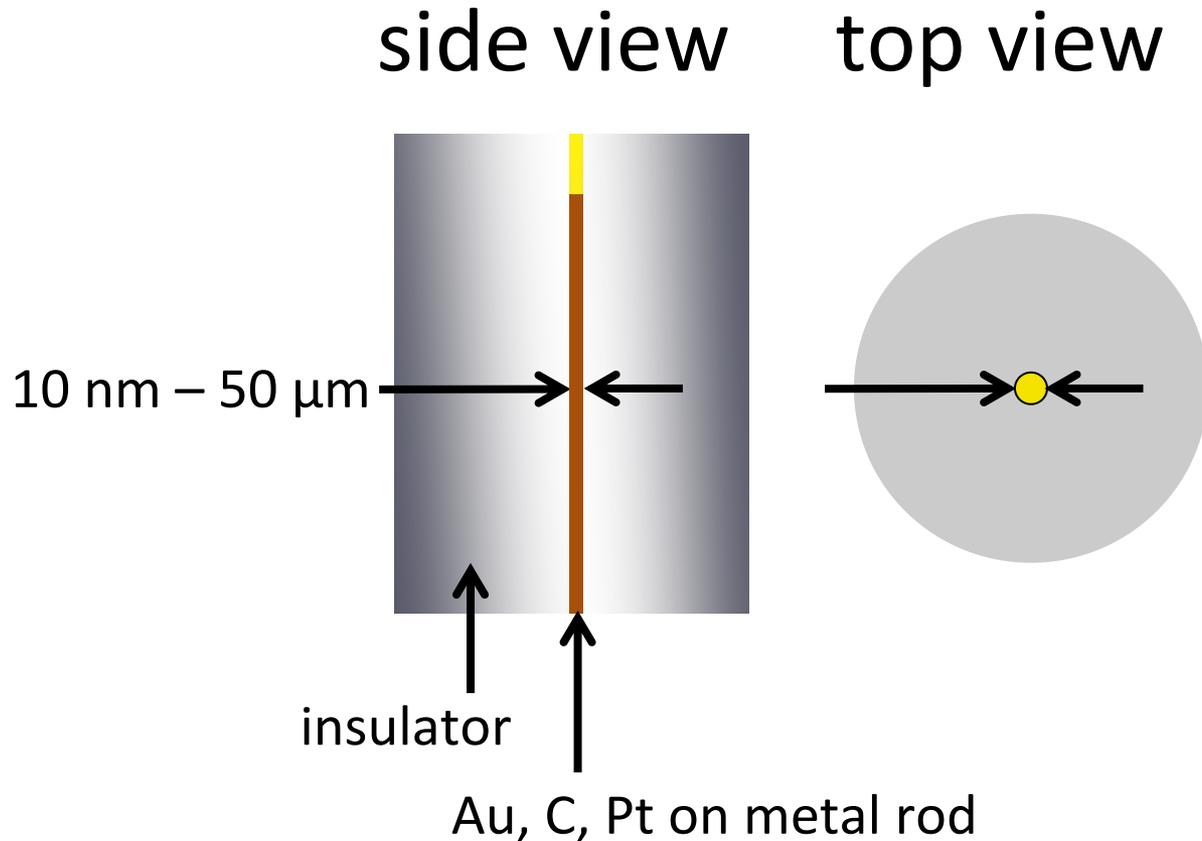


**How is current affected,
relative to the Cottrell
prediction?**

Let's design an experiment in which we intentionally operate in this radial diffusion limit the "entire" time!

RECALL:

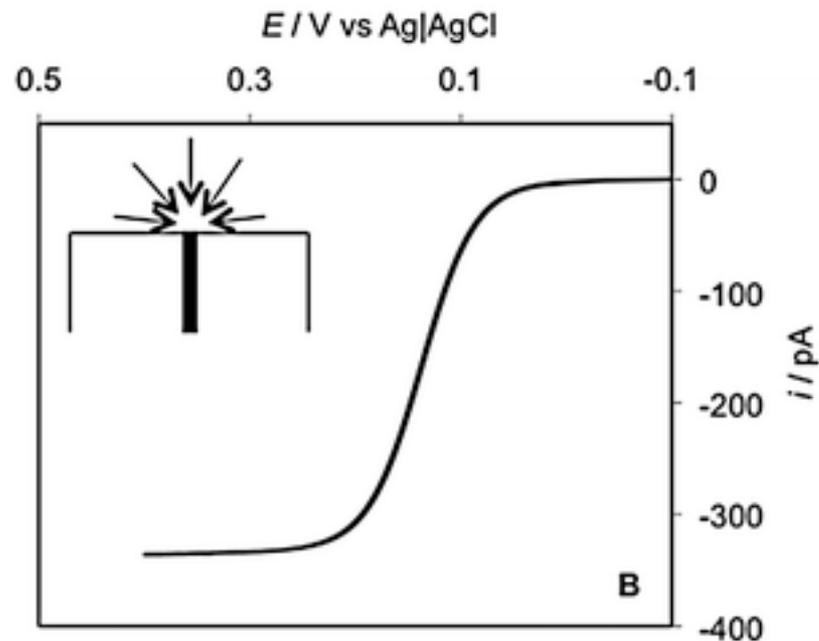
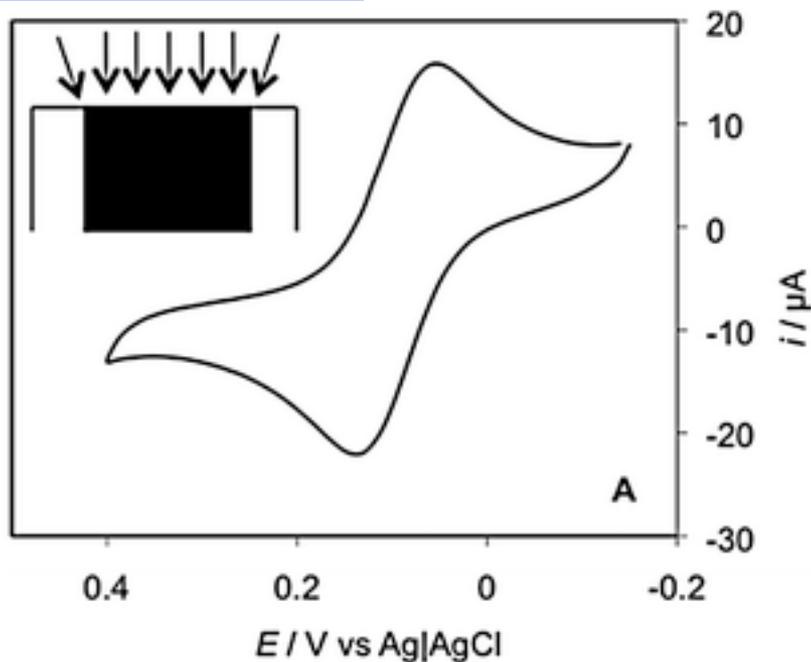
... well we actually start in the linear regime, and then switch over quickly...



... called "ultramicroelectrodes" or "UMEs"

RECALL:

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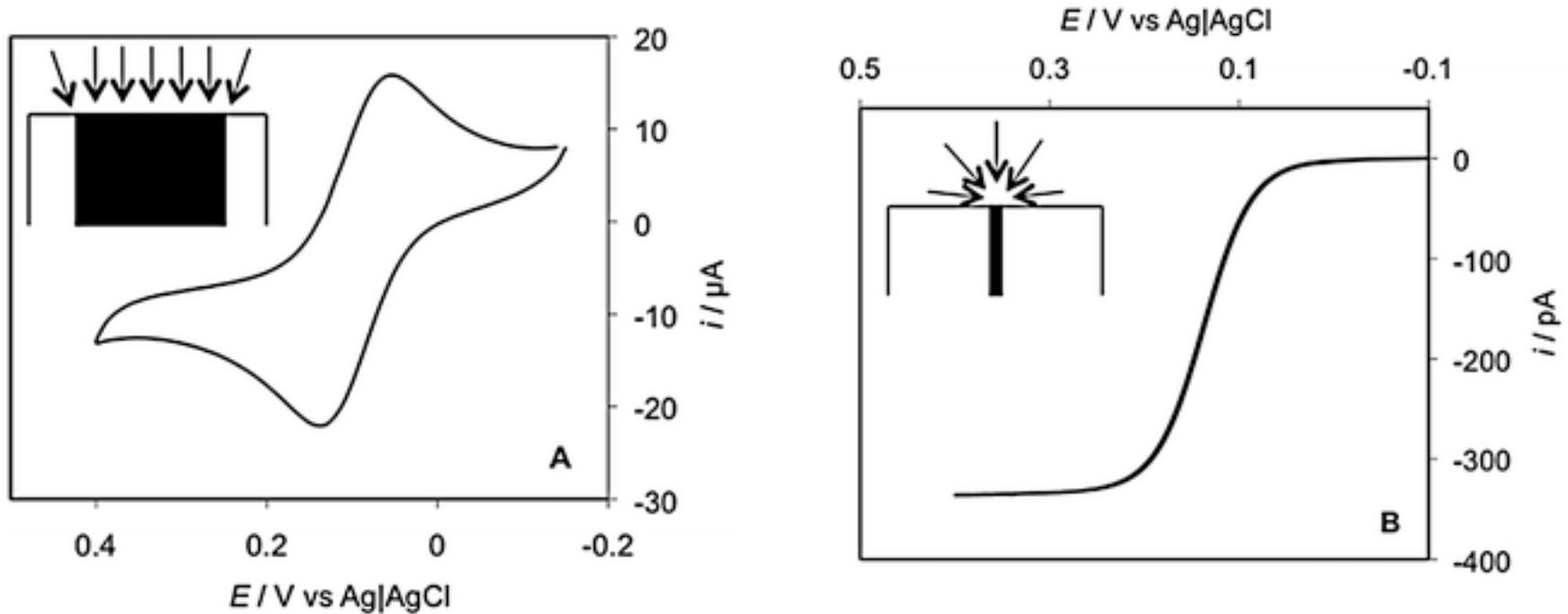
... steady-state is “often” reached at each applied potential at a UME during a sweep

... Let's quantify it... steady-state occurs when $v \ll RTD/(nFr_0^2)$

... v (mV s^{-1}) $\ll 26 \text{ mV} \times (D/r_0^2)$... for a *BASi* UME with $r_0 = 5 \mu\text{m}$...

... $26 \text{ mV} \times ((0.5 \times 10^{-5} \text{ cm}^2 \text{ s}^{-1}) / (0.5 \times 10^{-3} \text{ cm})^2) = 26 \text{ mV} \times (20 \text{ s}^{-1})$

... $v \ll 0.5 \text{ V s}^{-1}$... **Wow, you can still scan quite fast!**



- (1) Change sign of currents, because B&F indicates that positive current is cathodic ... this is likely because electrons are the charge carrier of current ... but physicists disagree; also, ions can be + or – so no need to focus on e^-
- (2) Mirror image the plot through the origin so that the axes look like math axes

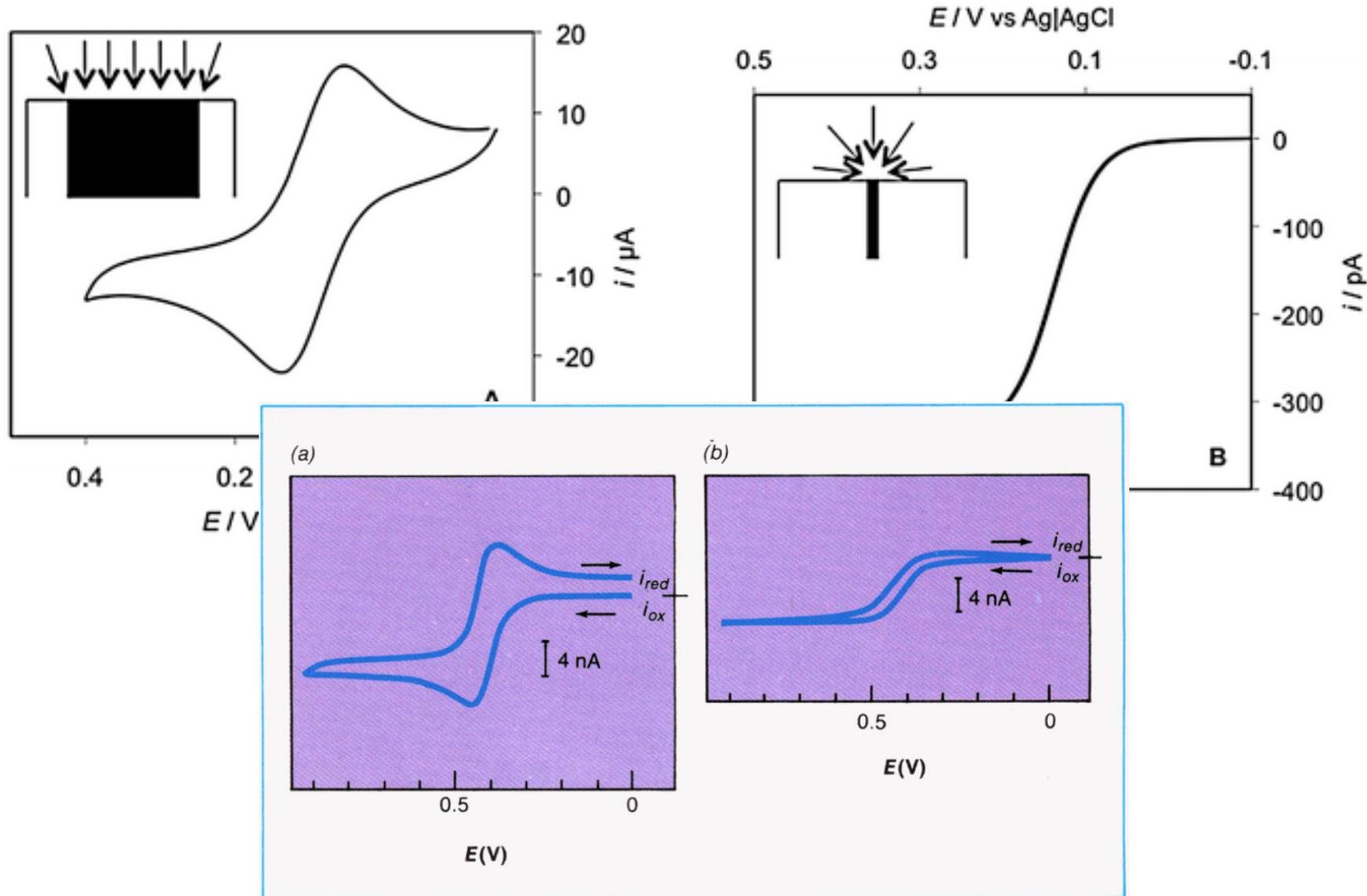
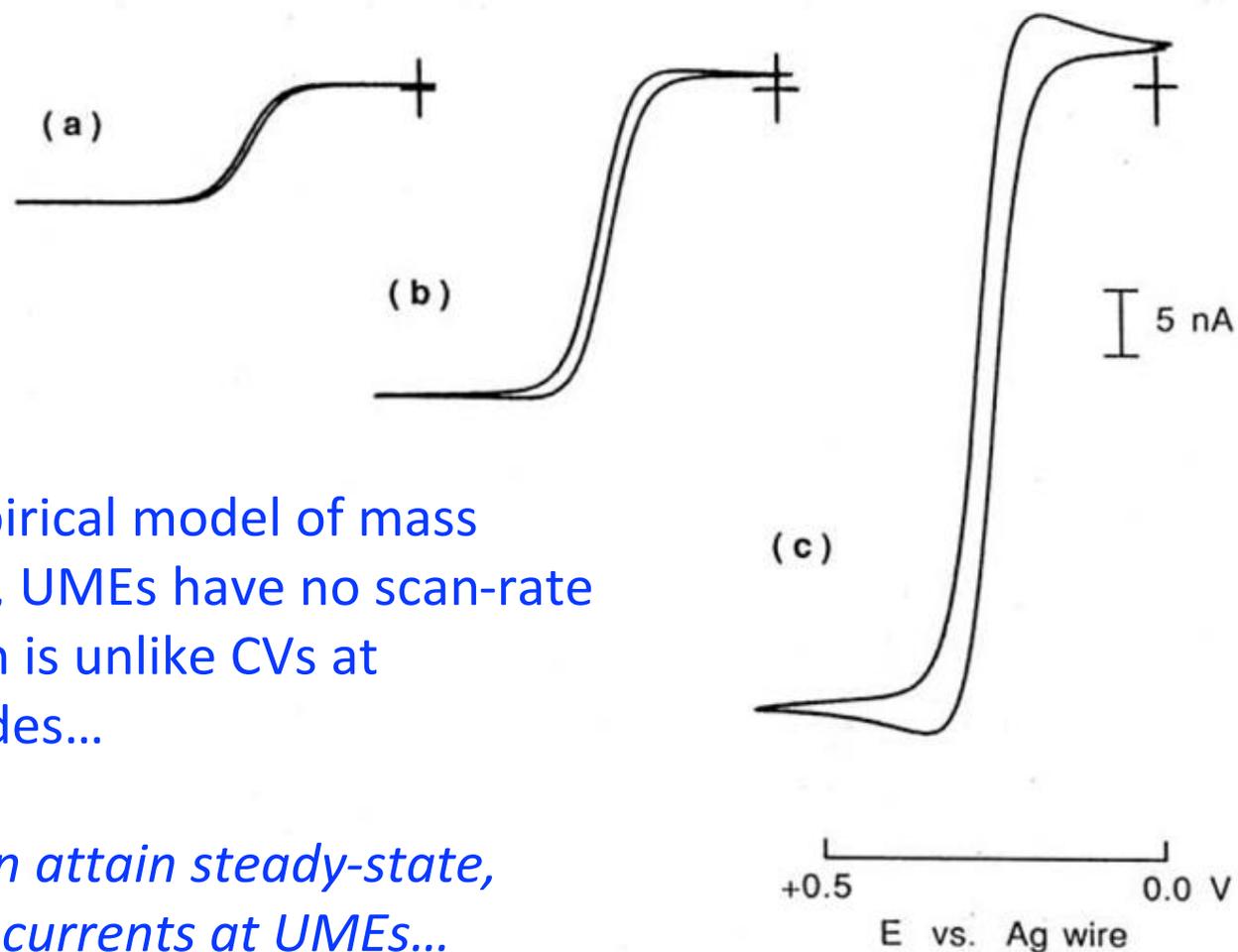


Figure 1. Cyclic voltammogram of ferrocene (1.0 mM) in acetonitrile with 0.1 M tetra-*n*-butylammonium perchlorate at a gold microdisk electrode ($r = 6.5 \mu\text{m}$). (a) 10 V s^{-1} scan rate. (b) 0.1 V s^{-1} scan rate



... akin to the semi-empirical model of mass transfer (Section 1.4.2), UMEs have no scan-rate dependence to I_p , which is unlike CVs at traditional disk electrodes...

... *this is because we can attain steady-state, mass-transport-limited currents at UMEs...*

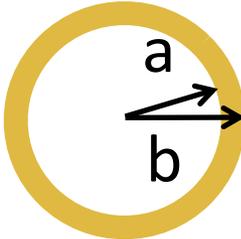
... ***without stirring at all!***

Figure 3. Cyclic voltammograms of 1 mM ferrocene in 0.1 M TBAPF₆/CH₃CN solutions obtained with (a) 10; (b) 25; (c) 50 μm -diameter Pt disks. The scan rate is 50 mV/s.

and the diffusion-limited current pre-factor depends on **electrode geometry**... 537

RECALL:

$$i_l = "x" n F D C^* r_0 \dots \text{but not scan rate!}$$

electrode geometry	"x"	
sphere	4π	
hemisphere	2π	
disk	4	
ring	$\frac{\pi^2 (b + a)}{r_0 \ln \left[16 \frac{(b + a)}{(b - a)} \right]}$	

... disk and ring-disk electrodes are real things that we will cover later

... how tiny (or large) is this diffusion-limited current (density)?

Let's assume $D = 0.5 \times 10^{-5} \text{ cm}^2 \text{ s}^{-1}$ and $C^* = 10^{-6} \text{ moles cm}^{-3}$ (1 mM)

r_0	$i_l = 4nFDC^*r_0$	$i_l/A = i_l/(\pi r_0^2)$
100 μm	20 nA	60 $\mu\text{A/cm}^2$
10 μm	2 nA	0.6 mA/cm ²
1 μm	0.2 nA	6 mA/cm ²
100 nm	20 pA	60 mA/cm ²
10 nm	2 pA	0.6 A/cm ²
1 nm	0.2 pA	6 A/cm ²

... super tiny currents...

... but huge current densities... = AWESOME!...

... for studies of electrocatalysis!

Again, Why do we care? One reason...

we need to push this up in order to perform meaningful measurements of the kinetics of rapid redox reactions

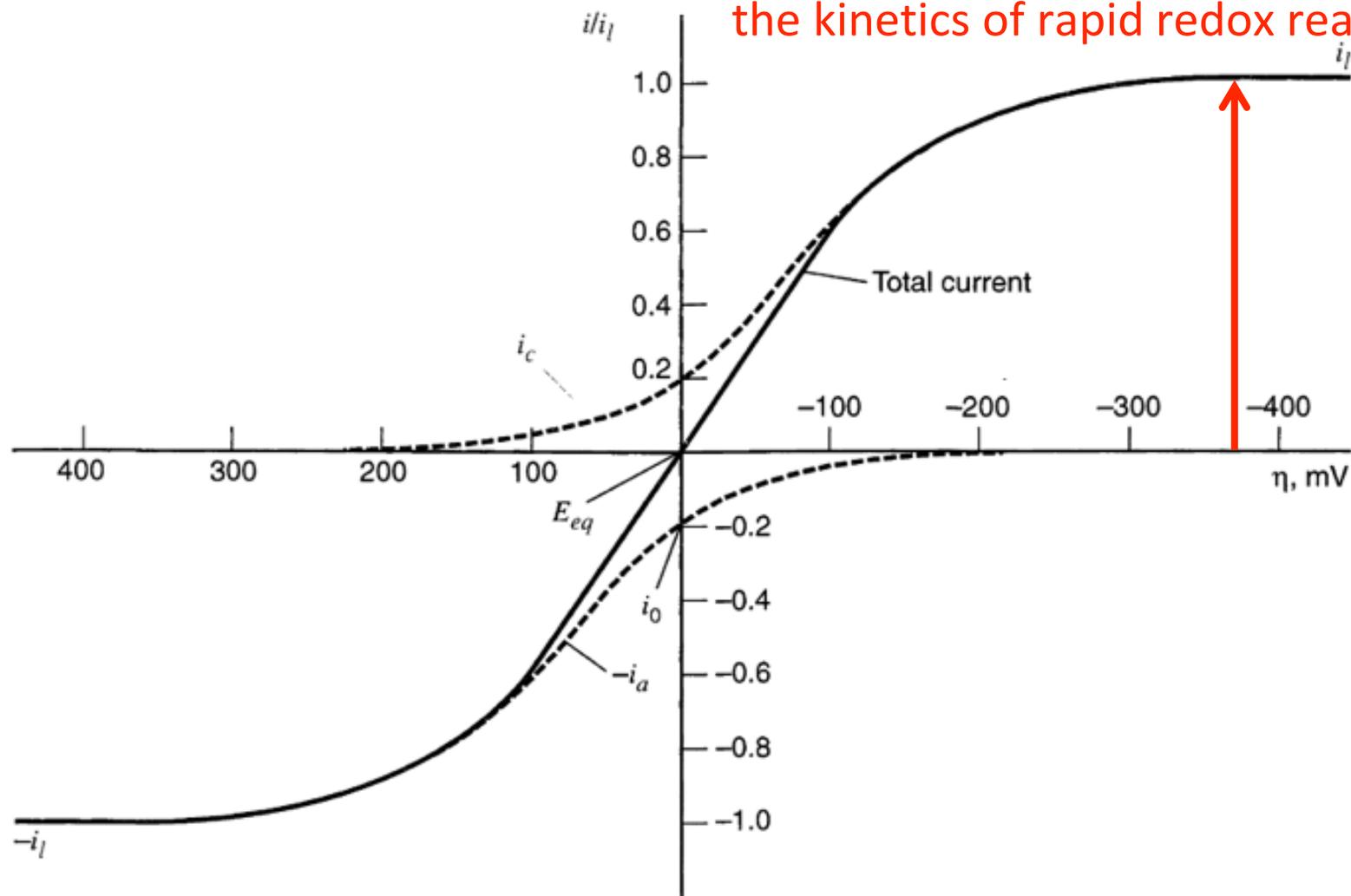


Figure 3.4.1 Current-overpotential curves for the system $O + e \rightleftharpoons R$ with $\alpha = 0.5$, $T = 298$ K, $i_{l,c} = -i_{l,a} = i_l$ and $i_0/i_l = 0.2$. The dashed lines show the component currents i_c and i_a .

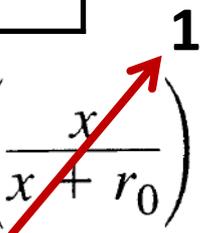
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r_0	$i_l = 4nFDC^*r_0$	$i_l/A = i_l/(\pi r_0^2)$
100 μm	20 nA	60 $\mu\text{A/cm}^2$
10 μm	2 nA	0.6 mA/cm ²
1 μm	0.2 nA	6 mA/cm ²
100 nm	20 pA	60 mA/cm ²
10 nm	2 pA	0.6 A/cm ²
1 nm	0.2 pA	6 A/cm ²

... so what is iR_u for these electrodes?

... well R_u is fairly large, but iR_u is nearly constant...

$$R_u = \frac{1}{4\pi\kappa r_0} \left(\frac{x}{x + r_0} \right)$$


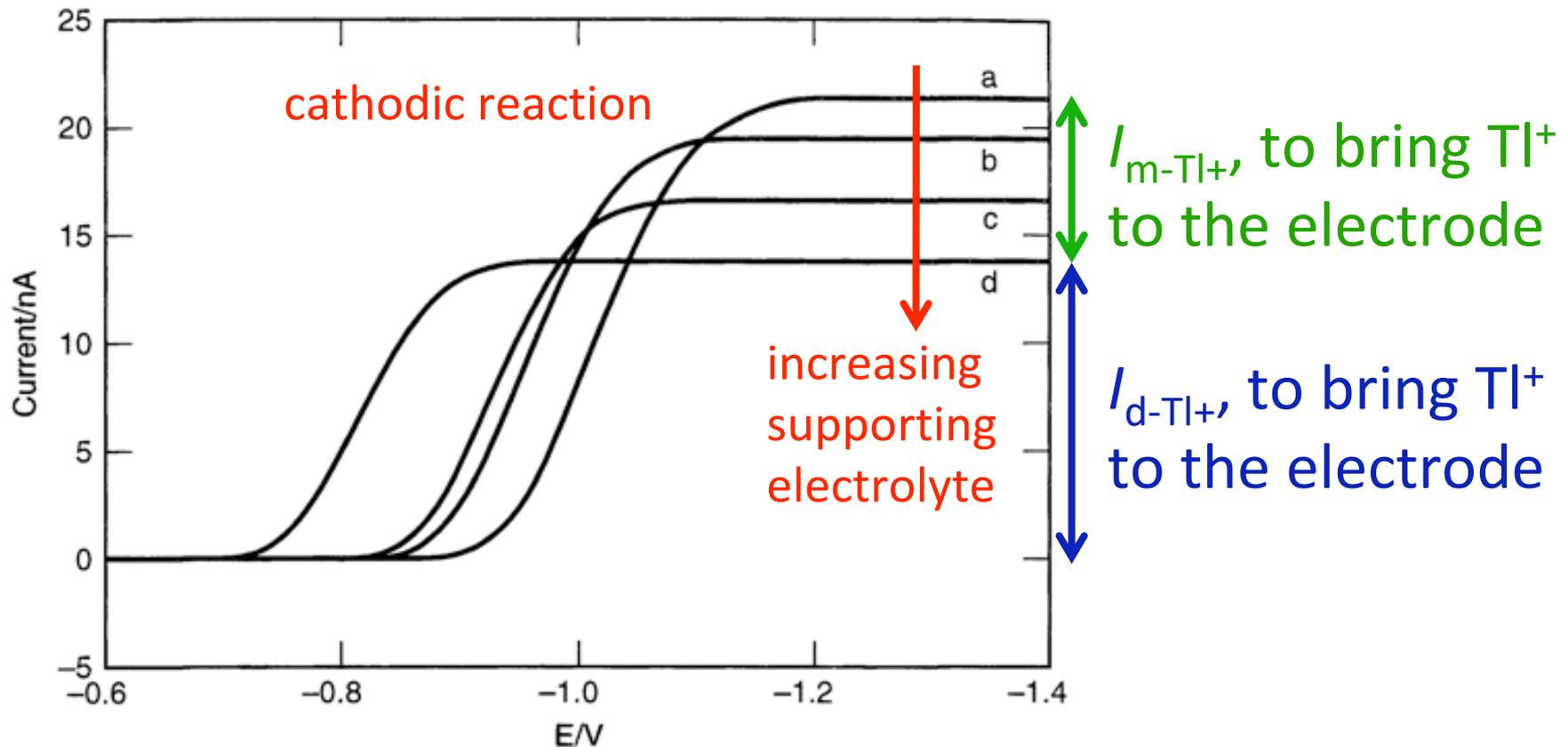


Figure 4.3.5 Voltammograms for reduction of 0.65 mM Tl_2SO_4 at a mercury film on a silver ultramicroelectrode (radius, 15 μm) in the presence of (a) 0, (b) 0.1, (c) 1, and (d) 100 mM $LiClO_4$. The potential was controlled vs. a Pt wire QRE whose potential was a function of solution composition. This variability is the basis for the shifts in wave position along the potential axis. [Reprinted with permission from M. Ciszowska and J. G. Osteryoung, *Anal. Chem.*, **67**, 1125 (1995). Copyright 1995, American Chemical Society.]

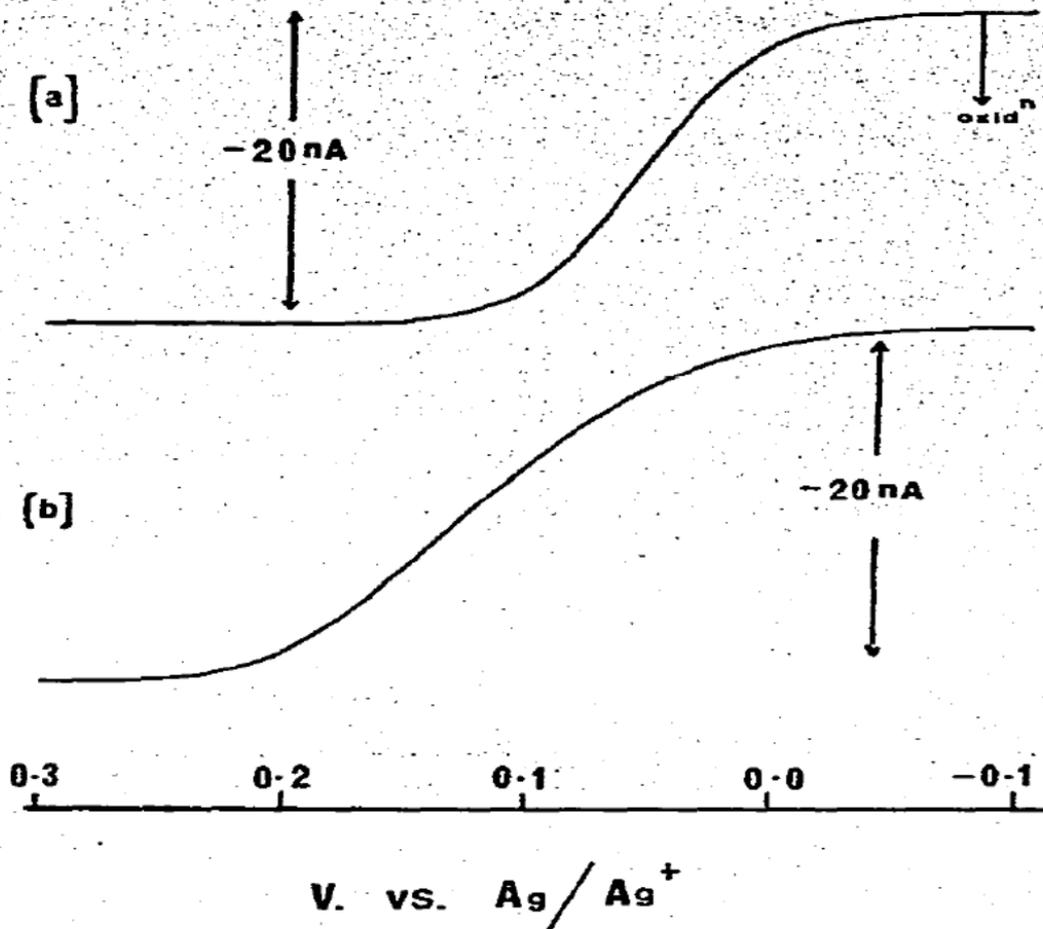


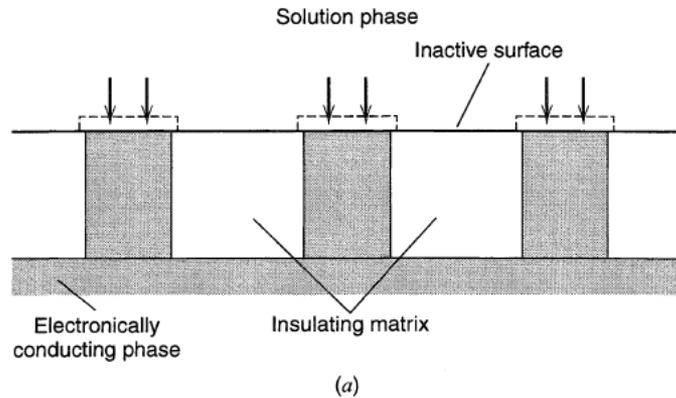
Fig. 2. Voltammogram for oxidation of 10^{-3} M ferrocene in acetonitrile at 18°C using a $25\text{-}\mu\text{m}$ radius Pt microelectrode in (a) presence and (b) absence of $0.1\text{ M Et}_4\text{NClO}_4$. Scan rate = 5 mV s^{-1} .

... additional/final points to address about UMEs:

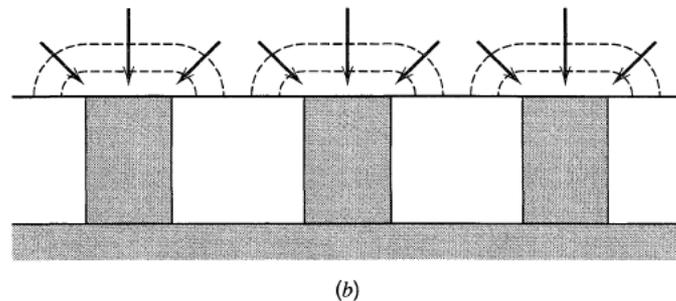
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- 1) You can buy them; how do you make them?
- 2) UME arrays and ensembles
- 3) Potential step experiments with UMEs...
- 4) How rapidly is steady-state attained?

... arrays of UMEs (from B&F)...



Semi-infinite linear diffusion



Semi-infinite mixed diffusion

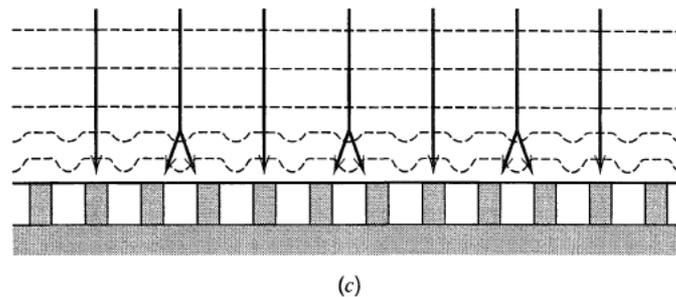


Figure 5.2.4 Evolution of the diffusion field during chronoamperometry at an electrode with active and inactive areas on its surface. In this case the electrode is a regular array such that the active areas are of equal size and spacing, but the same principles apply for irregular arrays. (a) Short electrolysis times, (b) intermediate times, (c) long times. Arrows indicate flux lines to the electrode.

Semi-infinite linear diffusion, *again*

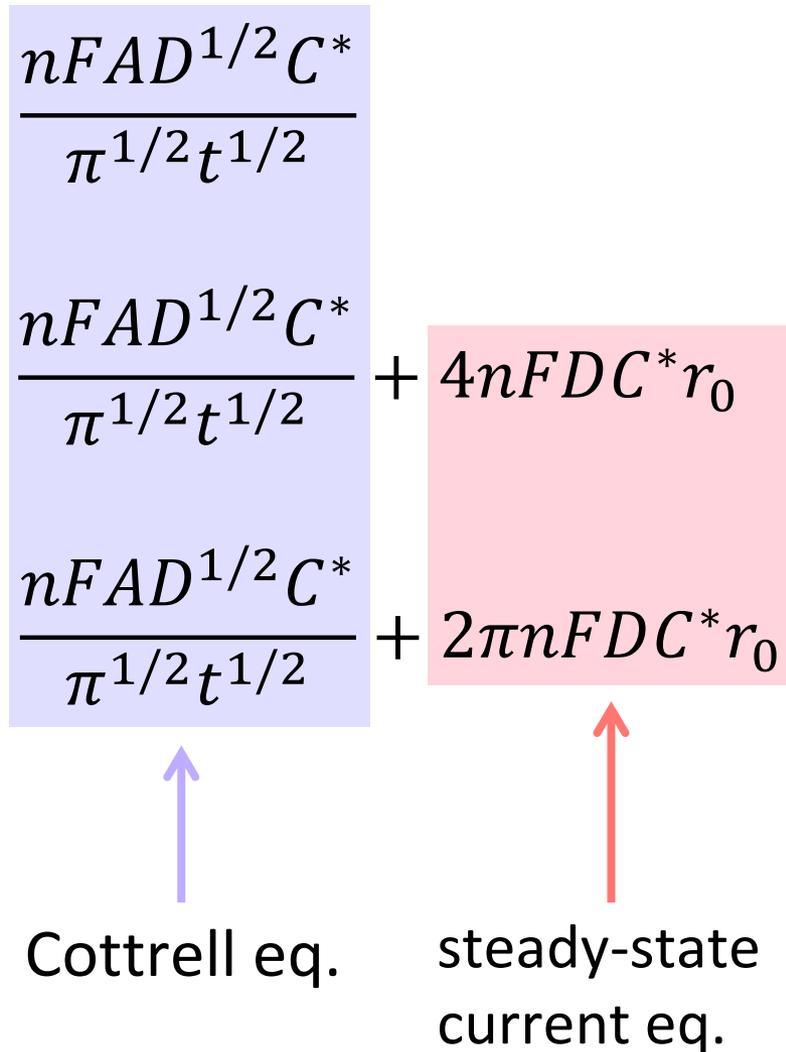
... no matter what some papers suggest... you can *never* beat overall linear diffusive flux over the total projected area of the diffusive solution region!

... this is because a UME must behave like a macroscopic electrode at sufficiently small times (and if neighboring electrodes are too close), right?... *Right!*

macroscopic planar electrode: $i(t) = \frac{nFAD^{1/2}C^*}{\pi^{1/2}t^{1/2}}$

disk UME: $i(t) = \frac{nFAD^{1/2}C^*}{\pi^{1/2}t^{1/2}} + 4nFDC^*r_0$

hemispherical UME: $i(t) = \frac{nFAD^{1/2}C^*}{\pi^{1/2}t^{1/2}} + 2\pi nFDC^*r_0$



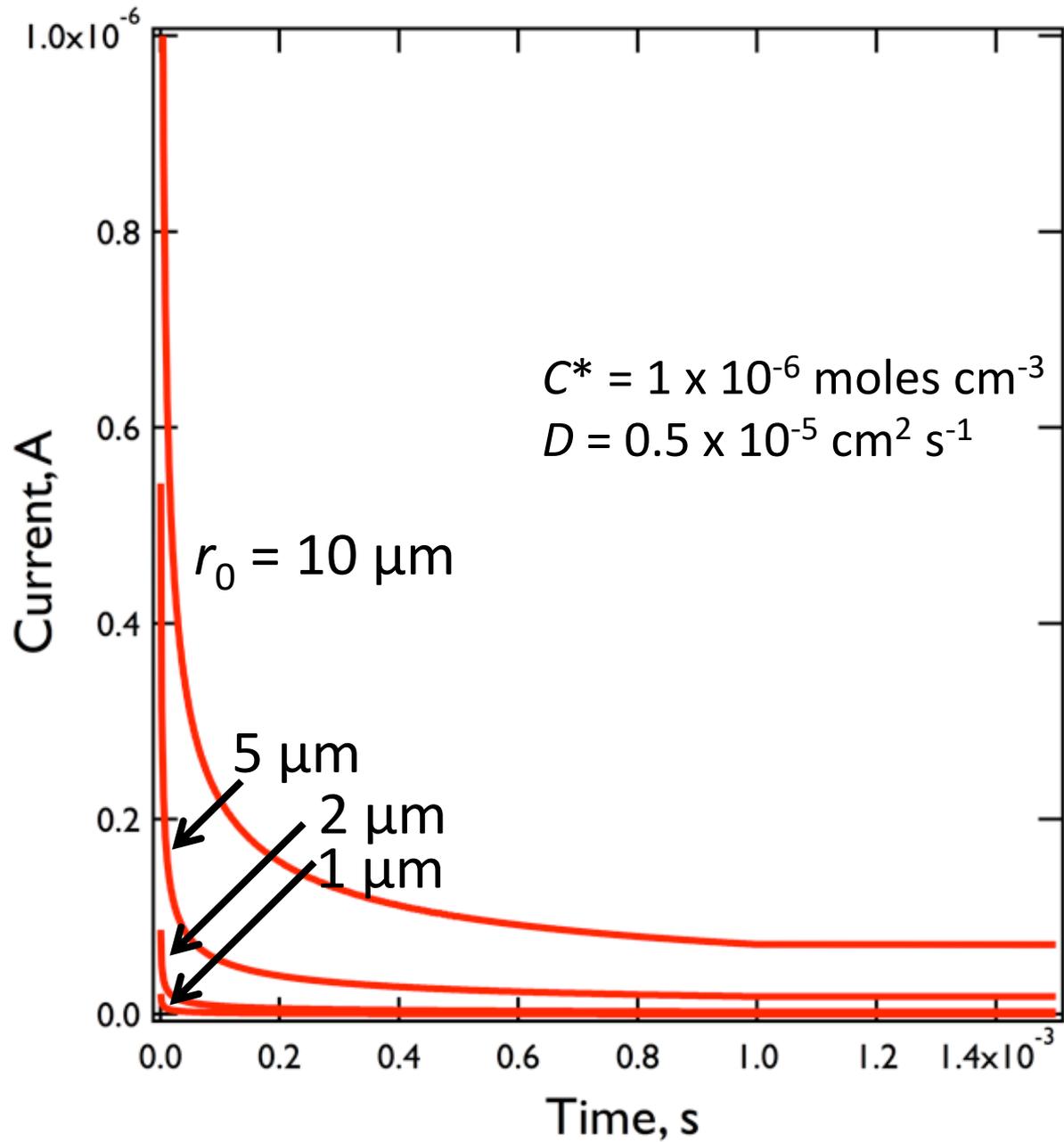
Cottrell eq. steady-state current eq.

... additional/final points to address about UMEs:

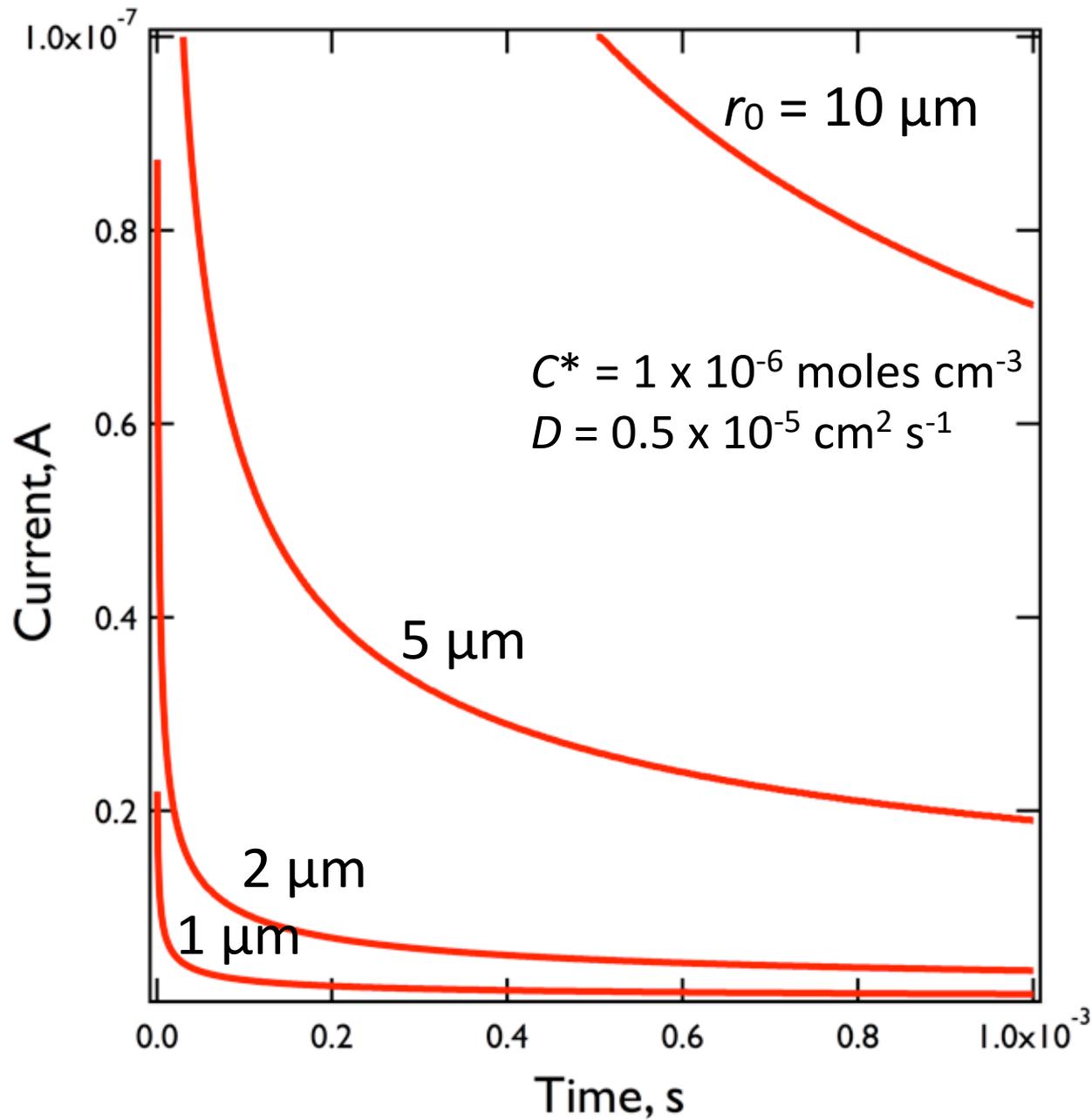
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- 1) You can buy them; how do you make them?
- 2) UME arrays and ensembles... **GOT IT!**
- 3) Potential step experiments with UMEs...
- 4) How rapidly is steady-state attained?

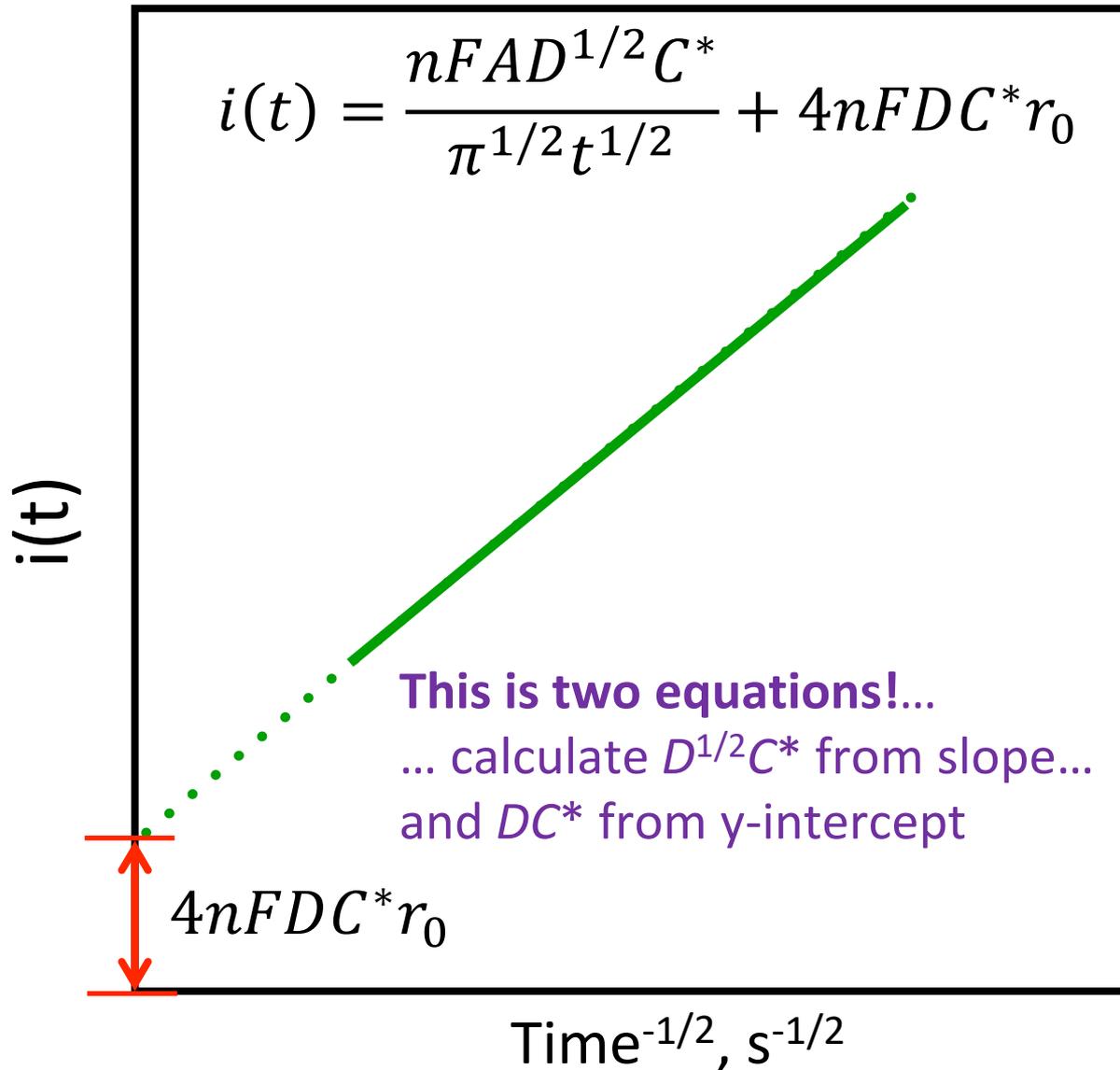
... at longer times, the current asymptotically approaches i_l ...



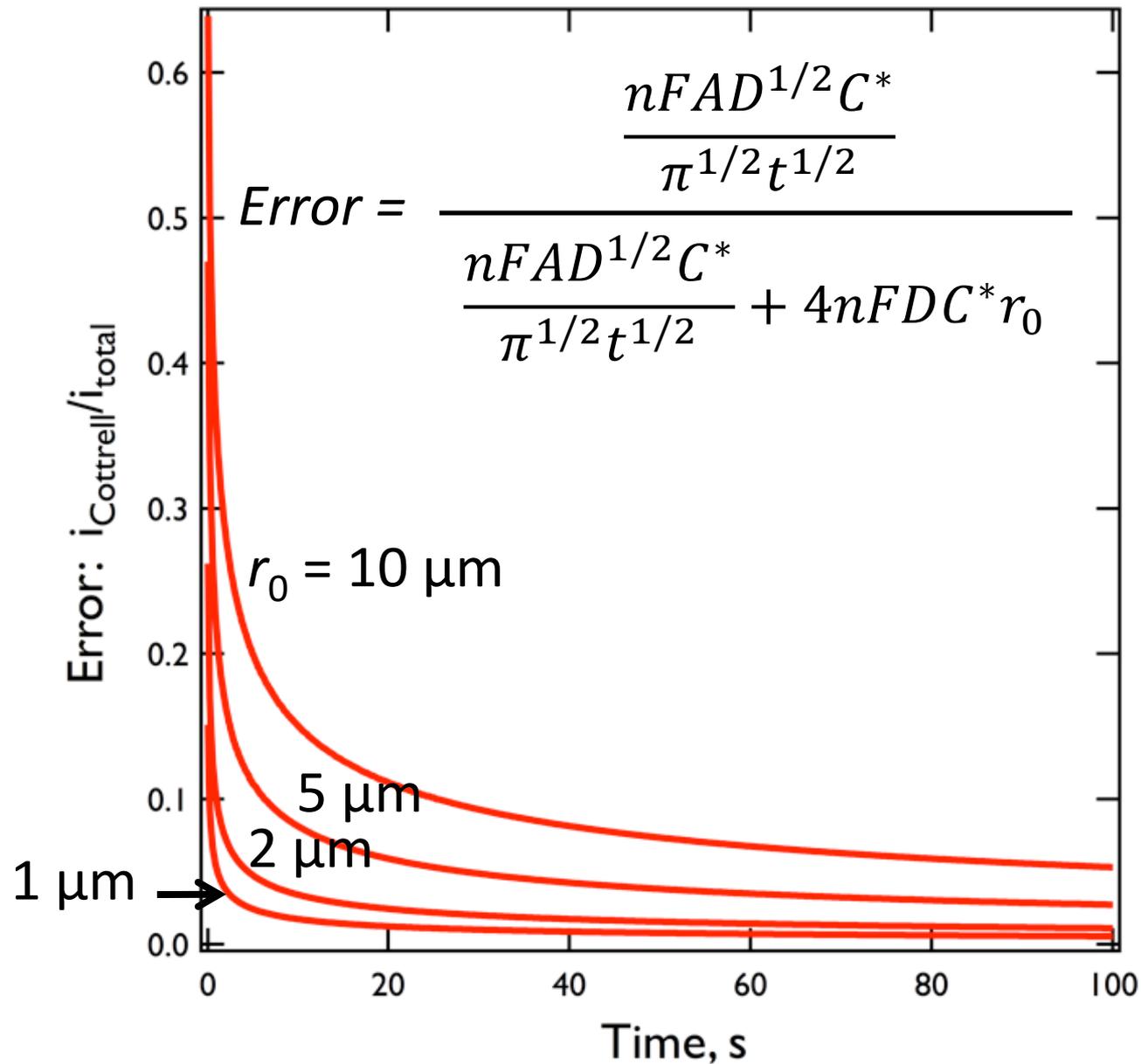
... at longer times, the current asymptotically approaches i_l ...



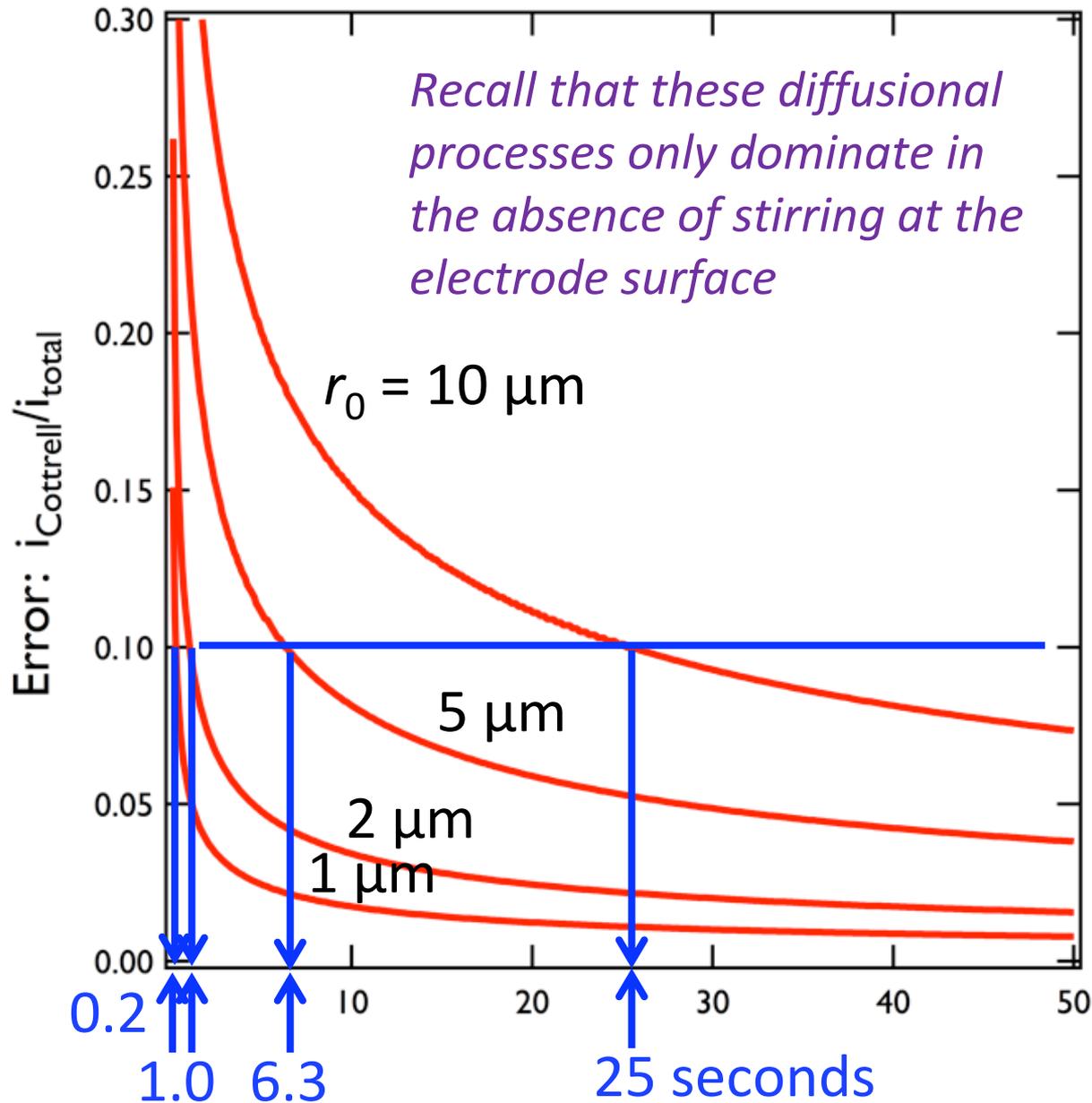
note that the $i(t)$ versus $(1/t^{1/2})$ plot no longer intersects "0" ...
 ... and D can be calculated without knowing C^* ... **How?**



... how long does it take for UME's to attain steady-state?



... how long does it take for UME's to attain steady-state?



the short time limit imposed by $R_u C_d$ for a macroscopic electrode is $\sim 100 \mu\text{s}$...

... but UMEs charge *much* faster!

r_0	r_0	$t_{\text{steady-state}}$	$R_u C_d^*$
10^{-3} cm	$10 \mu\text{m}$	25 s	$1.7 \mu\text{s}$
10^{-4} cm	$1 \mu\text{m}$	0.25 s	170 ns
10^{-5} cm	100 nm	2.5 ms	17 ns
10^{-6} cm	10 nm	$25 \mu\text{s}$	1.7 ns

$$\propto r_0^2$$

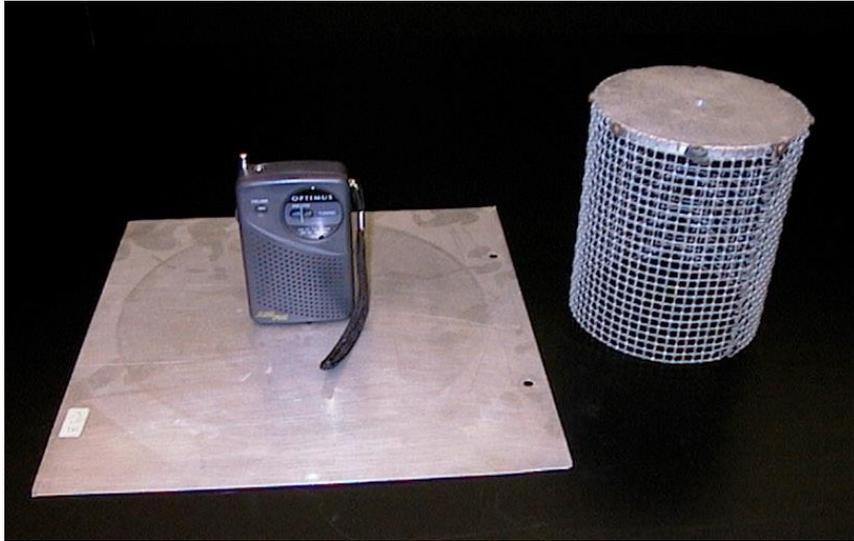
$$\propto r_0$$

$$\propto r_0^{-1}$$

... recall that R_u is nearly independent of the WE-RE separation:

$$R_u = \frac{1}{4\pi\kappa r_0} \left(\frac{x}{x+r_0} \right)$$

needed for $i_i < 10 \text{ nA}$: a Faraday cage...



... and a Keithley 428 programmable current amplifier grounded to the cage

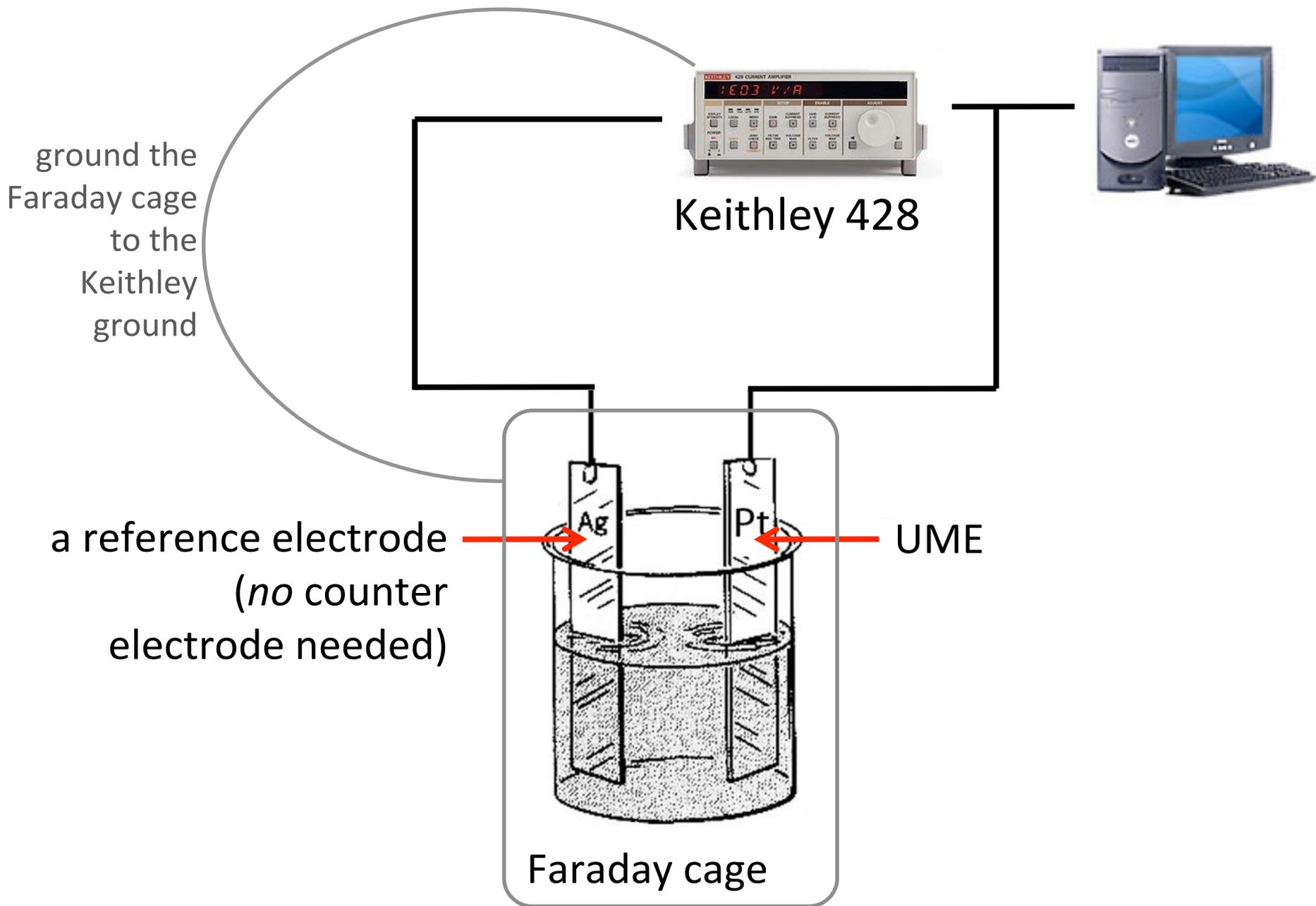


Key Features and Benefits:

- ▶ 2 μ s rise time
- ▶ 1.2fA rms noise
- ▶ Up to 10¹¹ V/A gain
- ▶ IEEE-488 interface

The Model 428-PROG Programmable Current Amplifier converts fast, small currents to a voltage, which can be easily digitized or displayed by an oscilloscope, waveform analyzer, or data acquisitionsystem. It uses a sophisticated "feedback current" circuit to achieve both fast rise times and sub-picoamp noise. The gain of the Model 428-PROG is adjustable in decade increments from 103V/A to 1011V/A, with selectable rise times from 2 μ s to 300ms.

... there are also a few things you will *not* be needing...



UME take-home messages:

After rapid double-layer charging...

... operate UMEs at either short times
= linear diffusion...

Can determine effects of rapid catalysis without complications from double-layer charging

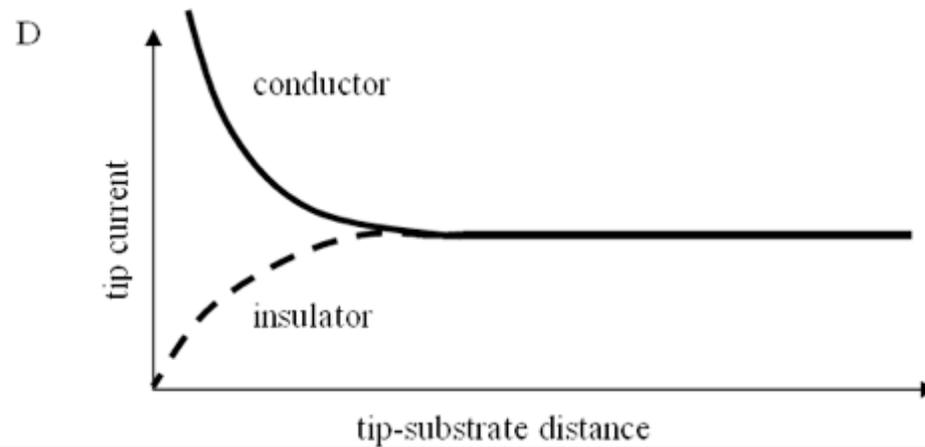
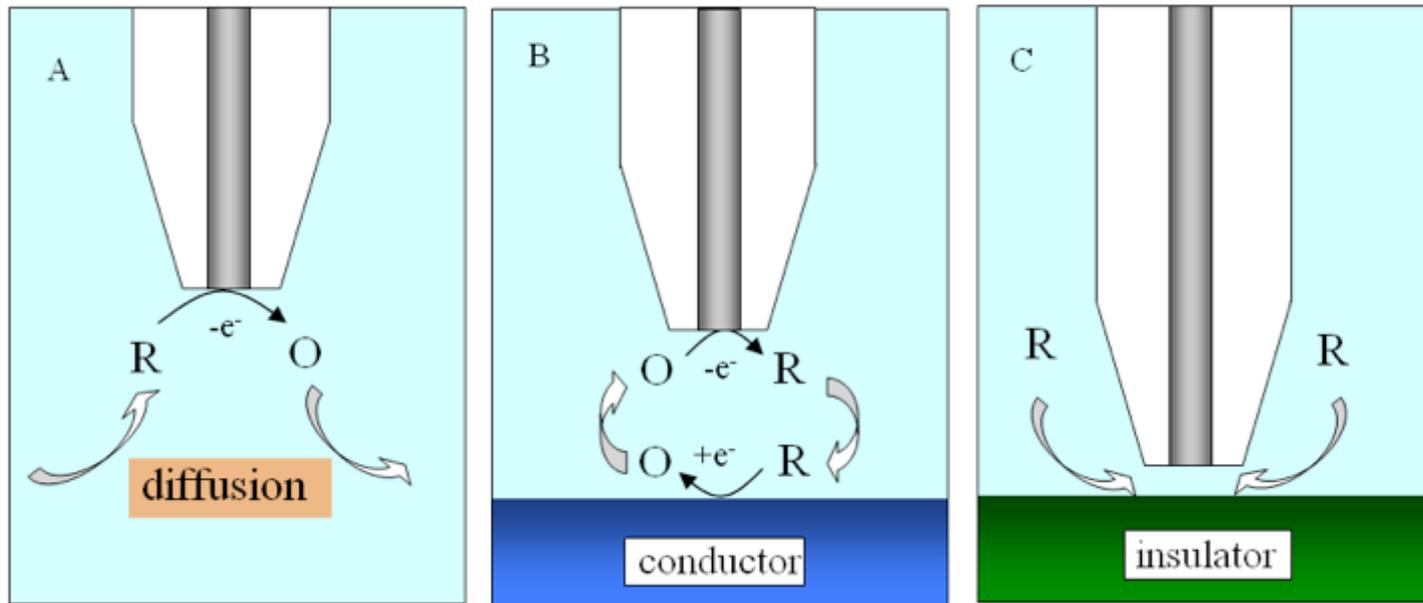
... or long times

= steady-state radial diffusion...

Can determine D without knowing C^*

We only have time to look at techniques/publications by one of these four UME pioneers... **Guess who?**

- ◆ Ralph “Buzz” Adams (d. Univ. Kansas)
- ◆ Mark Wightman (UNC Chapel Hill)
- ◆ Allen Bard (UT Austin)
- ◆ Henry White (Univ. Utah)



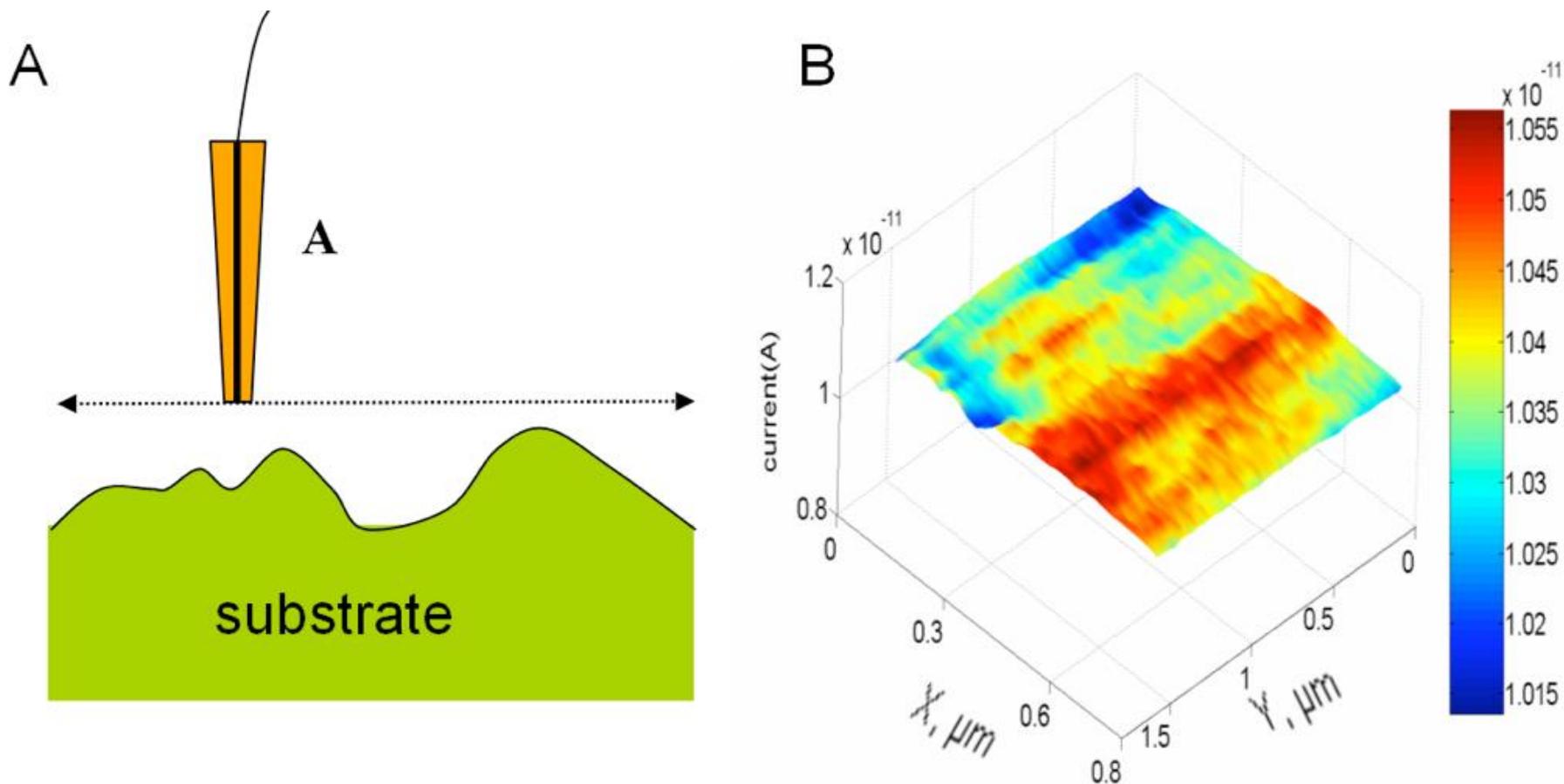
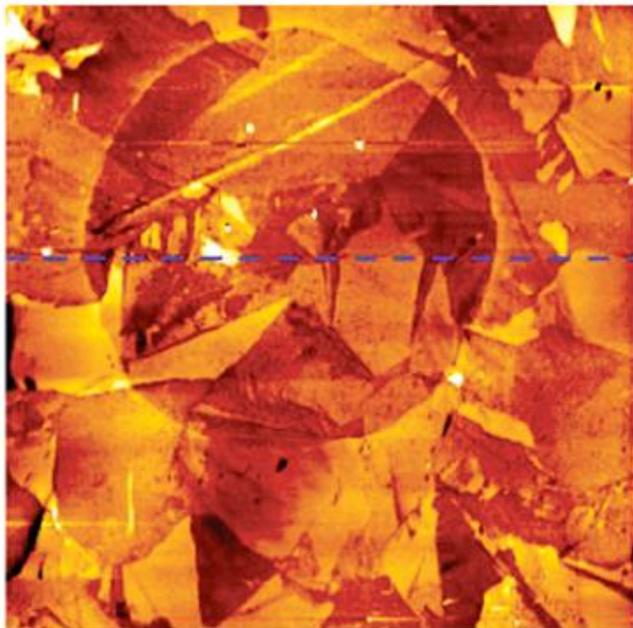


Fig. 5. Substrate imaging. (A) Tip is maintained at constant height while scanning horizontally. (B) SECM image of a portion ($1\ \mu\text{m} \times 1\ \mu\text{m}$) of a human breast cell membrane acquired with a $47\ \text{nm}$ radius ultramicroelectrode.

A



B

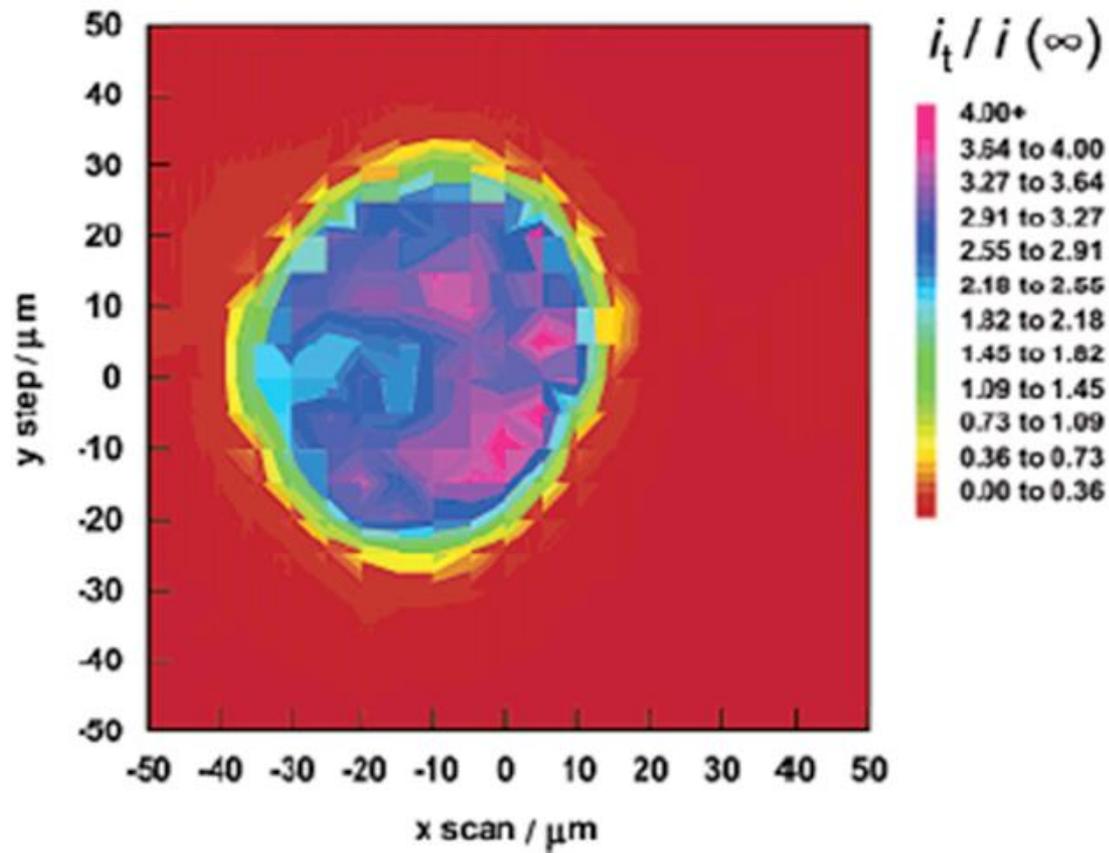


Fig. 7. (A) $80\ \mu\text{m} \times 80\ \mu\text{m}$ atomic force microscopy image of a disk-like region of boron-doped diamond. (B) SECM image over such a region.

Electrochemical Detection of Single Molecules

Fu-Ren F. Fan and Allen J. Bard*

The electrochemical behavior of a single molecule can be observed by trapping a small volume of a dilute solution of the electroactive species between an ultramicroelectrode tip with a diameter of ~ 15 nanometers and a conductive substrate. A scanning electrochemical microscope was used to adjust the tip-substrate distance (~ 10 nanometers), and the oxidation of [(trimethylammonio)methyl] ferrocene ($\text{Cp}_2\text{FeTMA}^+$) to $\text{Cp}_2\text{FeTMA}^{2+}$ was carried out. The response was stochastic, and anodic current peaks were observed as the molecule moved into and out of the electrode-substrate gap. Similar experiments were performed with a solution containing two redox species, ferrocene carboxylate ($\text{Cp}_2\text{FeCOO}^-$) and $\text{Os}(\text{bpy})_3^{2+}$ (bpy is 2,2'-bipyridyl).

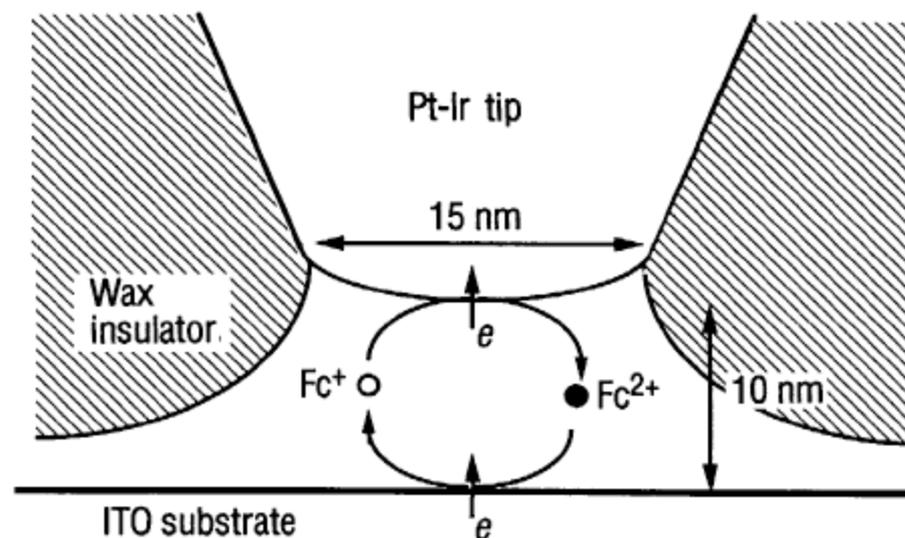
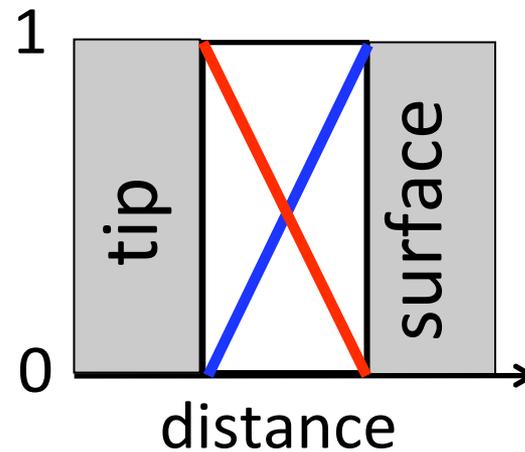
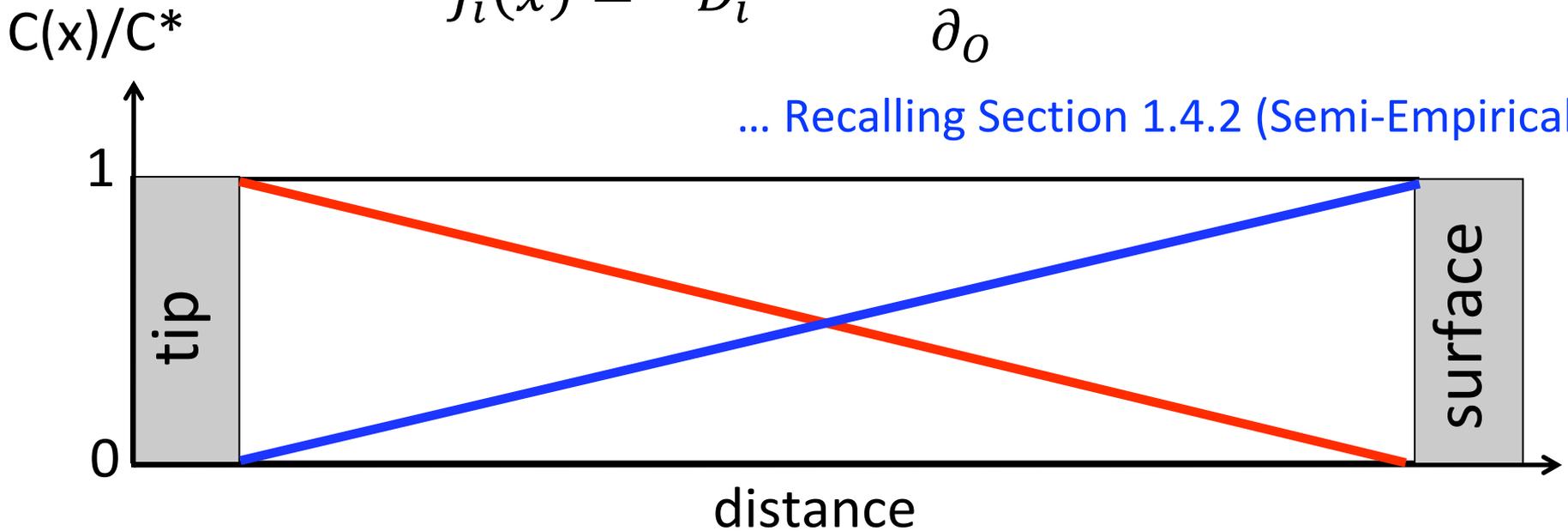


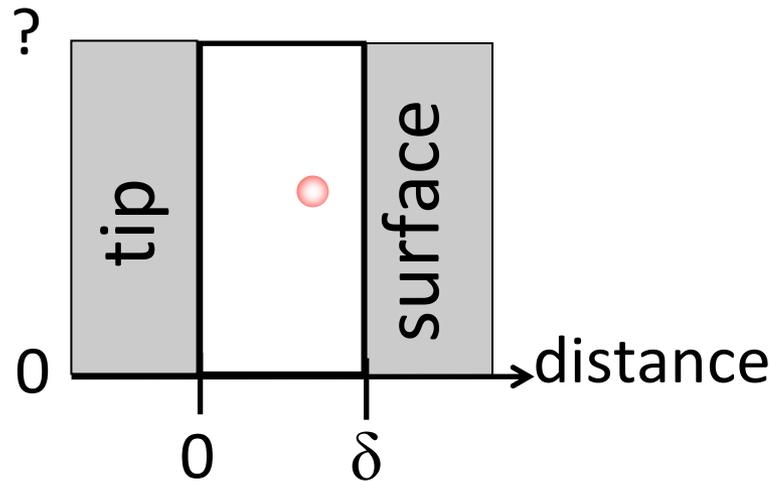
Fig. 1. Idealized schematic illustration of the tip geometry and the tip-substrate configuration used.

$$J_i(x) = -D_i \frac{C_o^* - C_o(x=0)}{\delta_o}$$

... Recalling Section 1.4.2 (Semi-Empirical)



A feasibility assessment... 1 molecule is trapped within a 1 μm x 1 μm x 10 nm volume between an SECM tip and a surface. What's the value of the limiting current?

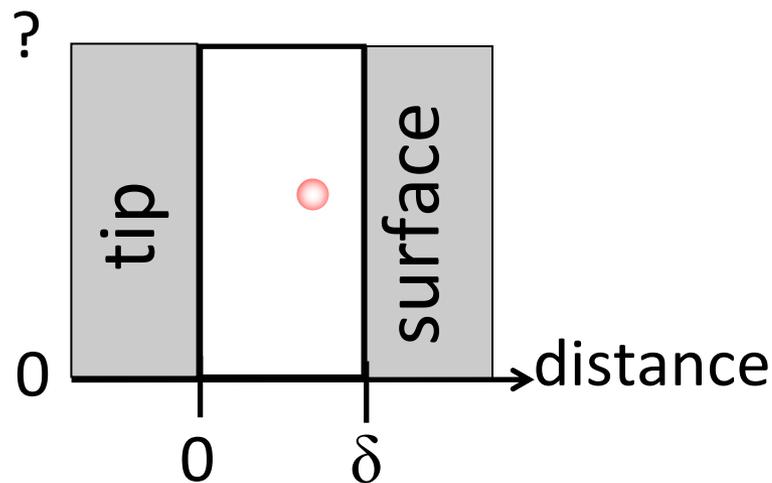


$$i = -nFAD \frac{\Delta C}{\delta} \quad m$$

$$C_{1 \text{ molecule}} = (1 \text{ molecule}) \left(\frac{1 \text{ mol}}{6.022 \times 10^{23} \text{ molecules}} \right) \left(\frac{1}{(10 \times 10^{-7} \text{ cm})(1 \times 10^{-4} \text{ cm})^2} \right)$$

$$C_{1 \text{ molecule}} = 1.66 \times 10^{-10} \text{ mol/cm}^3$$

A feasibility assessment... 1 molecule is trapped within a 1 μm x 1 μm x 10 nm volume between an SECM tip and a surface. What's the value of the limiting current?



$$i = -nFAD \frac{\Delta C}{\delta} \quad m$$

$$i = (1 \text{ eq/mol})(96485 \text{ C/eq})(1 \times 10^{-4} \text{ cm})^2 (1 \times 10^{-5} \text{ cm}^2/\text{s}) \left(\frac{1.66 \times 10^{-10} \text{ mol/cm}^3}{10 \times 10^{-7} \text{ cm}} \right)$$

$$= 1.6 \times 10^{-12} \text{ A} = 1.6 \text{ pA}$$

... so we're talking about pA's. We can measure that!

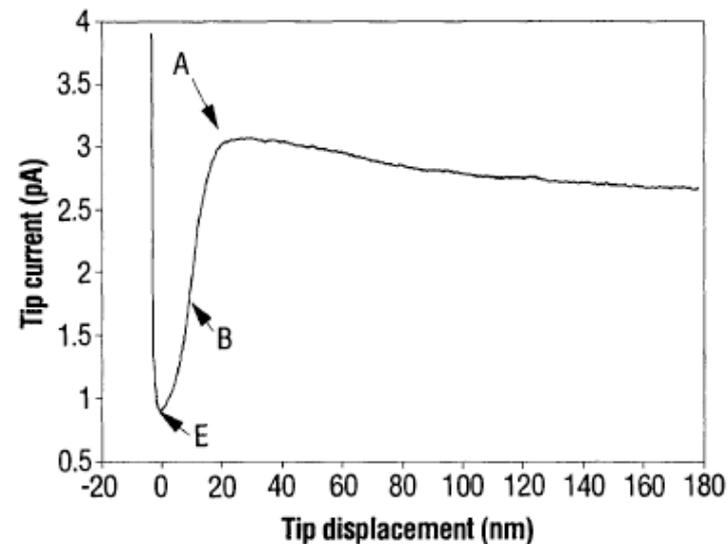
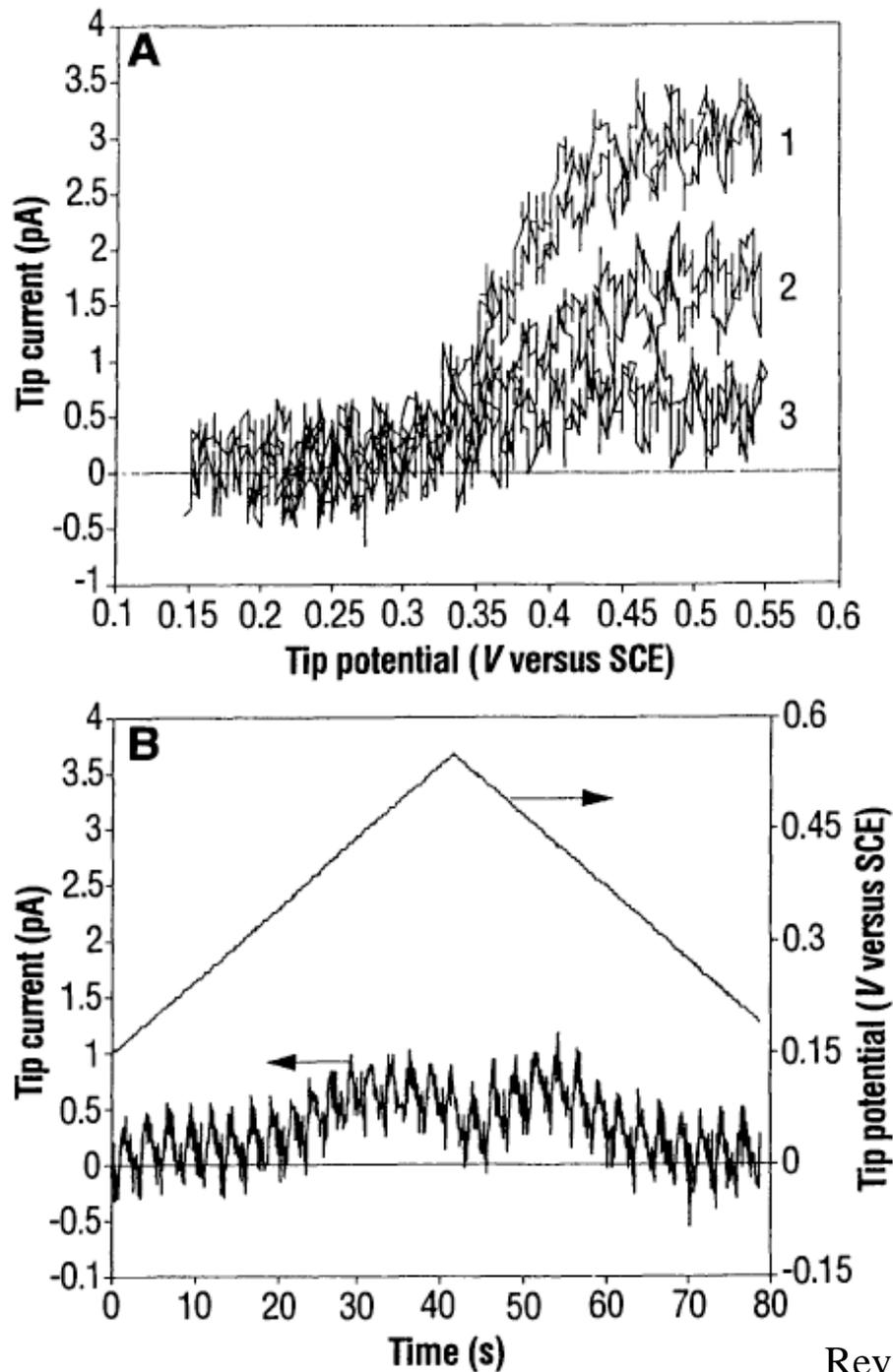


Fig. 2. Dependence of tip current on relative tip displacement over a conductive ITO substrate in a solution containing 2 mM $\text{Cp}_2\text{FeTMA}^+$ and 2.0 M NaNO_3 . The ITO substrate was biased at -0.3 V versus SCE and the tip was biased at 0.55 V, where the redox reactions on both electrodes were diffusion-controlled. The tip moved to the substrate surface at a rate of 30 \AA/s . The data were 300 points averaged at each location. The significance of points A, B, and E is discussed in the text. See also Fig. 4.

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- Section 5.1: Overview of step experiments
- Section 5.2: Potential step under diffusion controlled
- Sections 5.3 & 5.9: Ultramicroelectrodes
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A detailed review of Section 4.4.2 and Chapter 5

- Fick's Second Law of Diffusion
- Linear Diffusion = time-dependent current (Cottrell Equation)
- Anson Plots for surface adsorbed species
- Radial Diffusion = time-independent current (at steady-state)
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