

# Lecture B4

## Vibrational Spectroscopy, Part 2

## Quantum Mechanical SHO.

$$\hat{H}|\Psi_n\rangle = E_n|\Psi_n\rangle$$

WHEN we solve the Schrödinger equation, we always obtain two things:

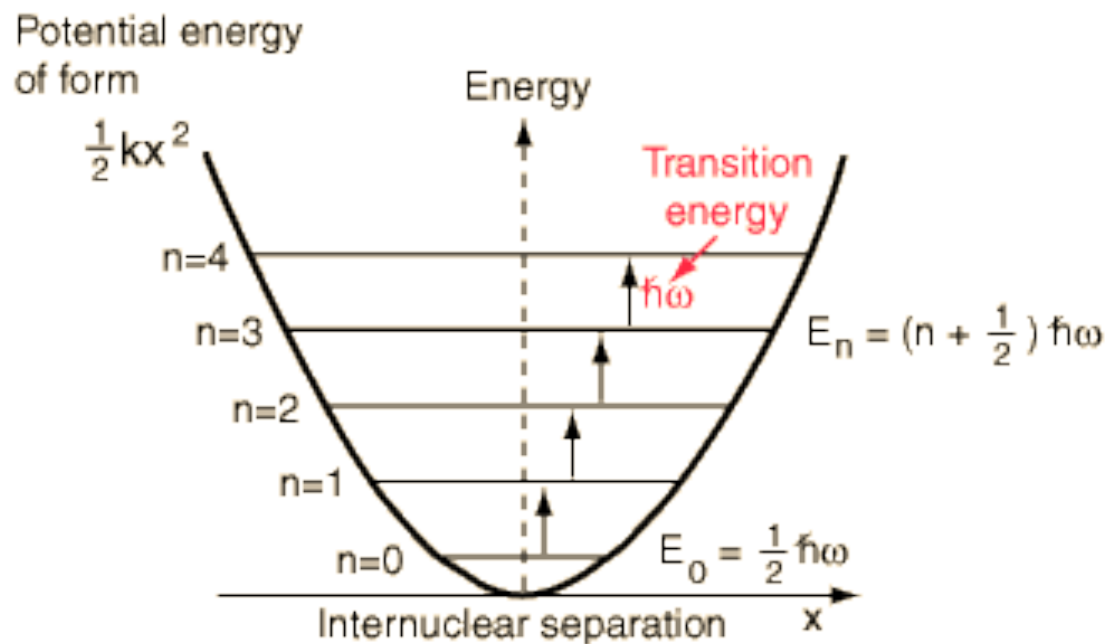
1. a set of eigenstates,  $|\Psi_n\rangle$ .
2. a set of eigenstate energies,  $E_n$

$$|\psi_n\rangle$$

$$n = 0, 1, 2, \dots$$

$$E_n = h\nu(n + \frac{1}{2})$$

QM predicts the existence of discrete, evenly spaced, vibrational energy levels for the SHO.



$$E_n = h\nu(n + \frac{1}{2})$$

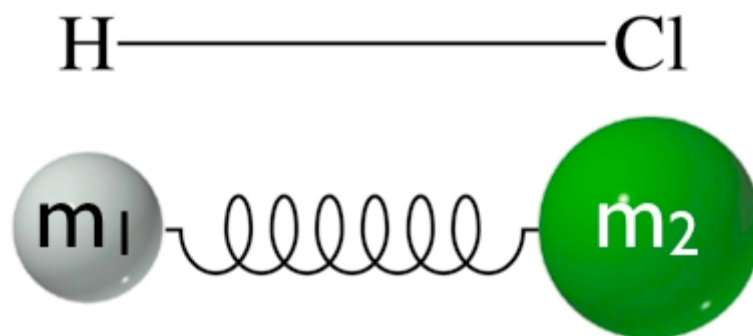
$$n = 0, 1, 2, 3, \dots$$

Notes:

- For the ground state ( $n=0$ ),  $E = \frac{1}{2}h\nu$ . This is called the zero point energy.
- Optical selection rule -- SHO can absorb or emit light with a  $\Delta n = \pm 1$

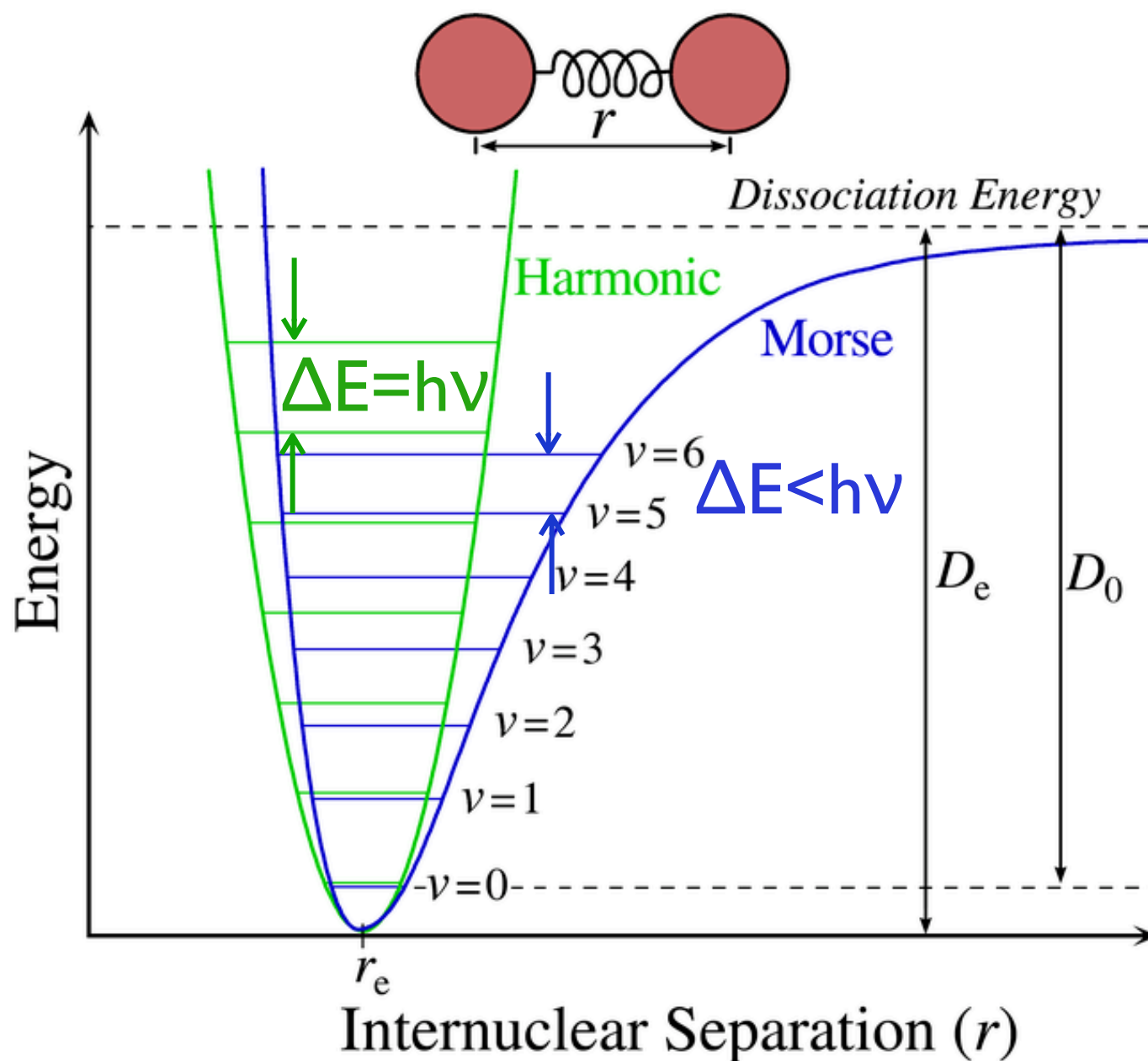
The IR absorption spectrum for a diatomic molecule, such as HCl:

diatomic

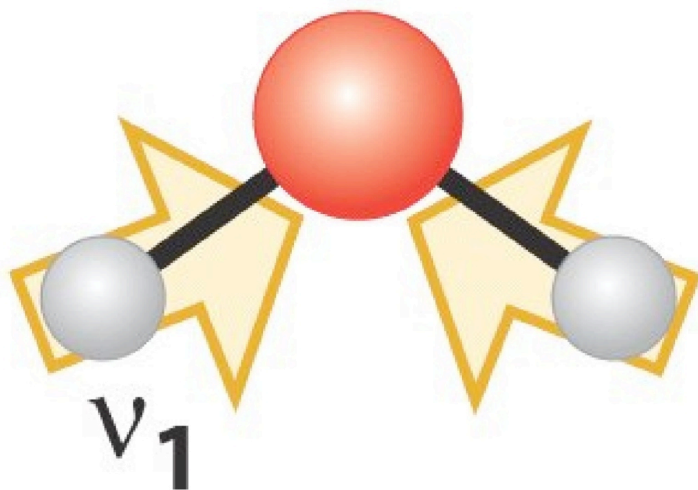


- Optical selection rule 1 -- SHO can absorb or emit light with a  $\Delta n = \pm 1$
- Optical selection rule 2 -- a change in molecular dipole moment ( $\Delta\mu/\Delta x$ ) must occur with the vibrational motion.  
(Note:  $\mu$  here means dipole moment).

If a more realistic Morse potential is used in the Schrödinger Equation, these energy levels get scrunched together...



The vibrational spectroscopy of polyatomic molecules gets more interesting...



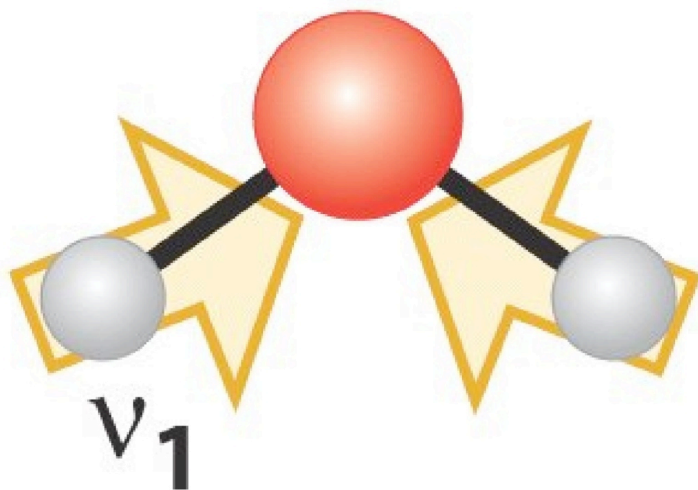
For diatomic or linear molecules:  $3N-5$  modes

For nonlinear molecules:  $3N-6$  modes



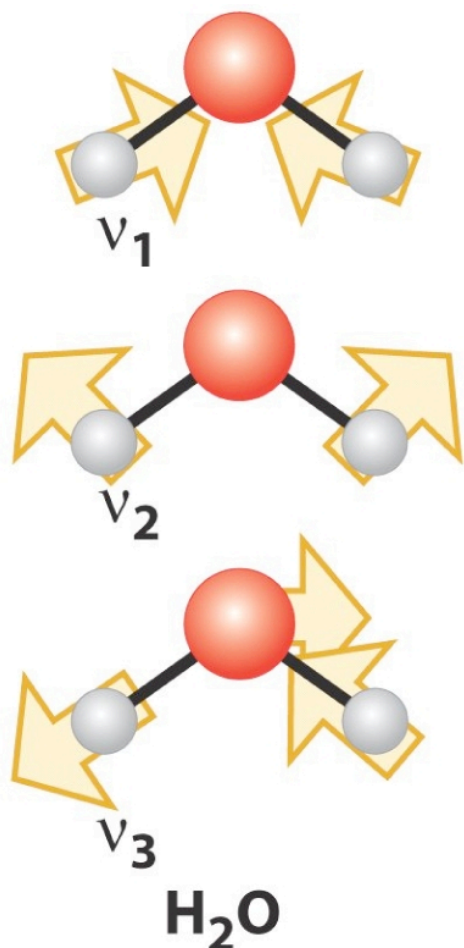
N = number of atoms in molecule

The vibrational spectroscopy of polyatomic molecules gets more interesting...



- Optical selection rule 1 -- SHO can absorb or emit light with a  $\Delta n = \pm 1$
- Optical selection rule 2 -- a change in molecular dipole moment ( $\Delta\mu/\Delta x$ ) must occur with the vibrational motion of a mode.

Consider H<sub>2</sub>O (a nonlinear molecule):

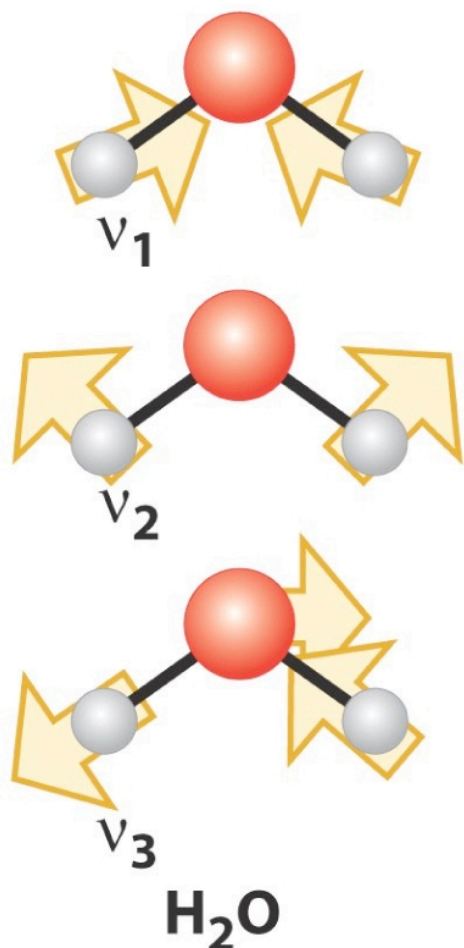


$$3N-6 = 3(3)-6 = 3$$

- Optical selection rule 2 -- a change in molecular dipole moment ( $\Delta\mu/\Delta x$ ) must occur with the vibrational motion of a mode.



Consider H<sub>2</sub>O (a nonlinear molecule):



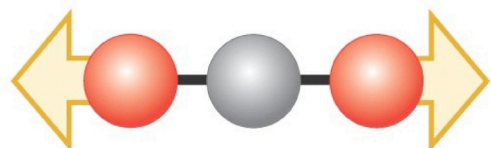
$$3N-6 = 3(3)-6 = 3$$

All bands are observed  
in the IR spectrum.



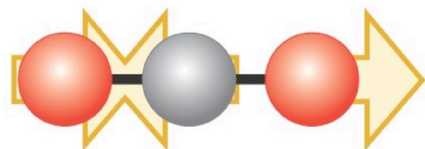
Optical selection rule 2 -- a change in molecular dipole moment ( $\Delta\mu/\Delta x$ ) must occur with the vibrational motion of a mode.

Consider CO<sub>2</sub> (a linear molecule):



$\nu_1$  Symmetrical stretch

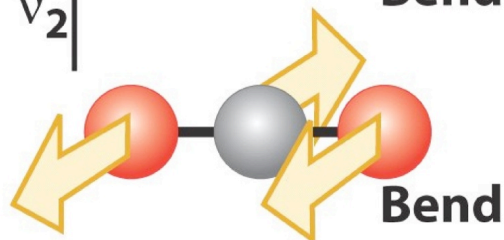
$$3N-5 = 3(3)-5 = 4$$



$\nu_3$  Antisymmetrical stretch



$\nu_2$  Bend



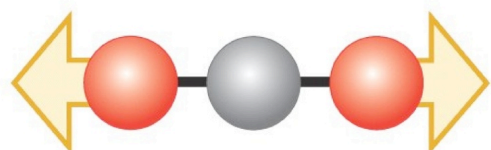
Bend

CO<sub>2</sub>



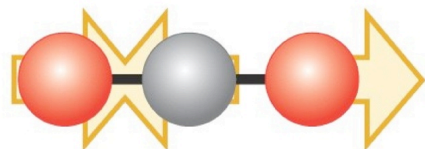
Optical selection rule 2 -- a change in molecular dipole moment ( $\Delta\mu/\Delta x$ ) must occur with the vibrational motion of a mode.

Consider CO<sub>2</sub> (a linear molecule):



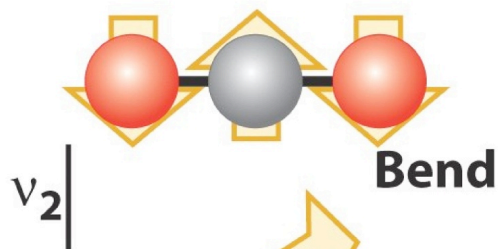
$\nu_1$  Symmetrical stretch

$$3N-5 = 3(3)-5 = 4$$



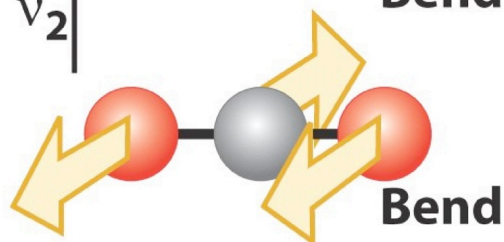
$\nu_3$  Antisymmetrical stretch

The motion of the symmetrical stretch ( $\nu_1$ ) does not create a new dipole moment, so it is not observed in the IR spectrum.



Bend

$\nu_2$



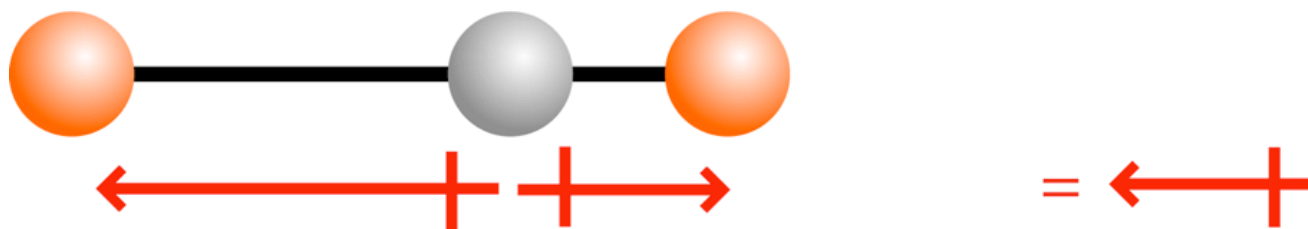
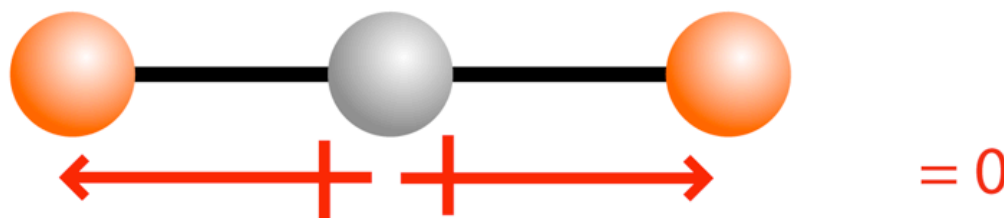
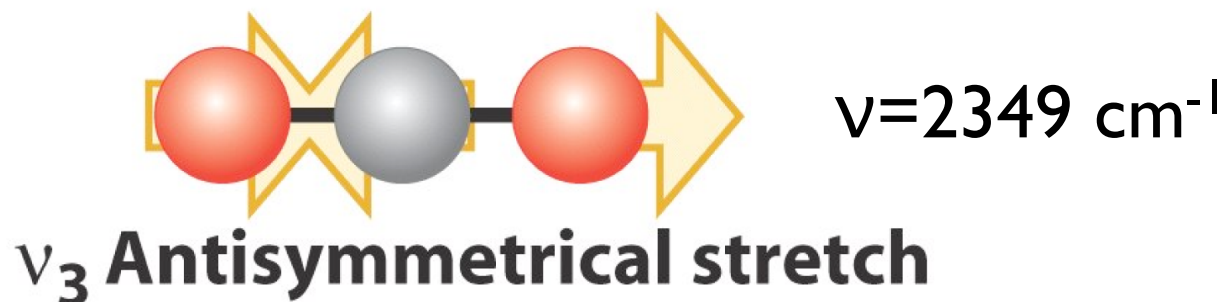
Bend

CO<sub>2</sub>

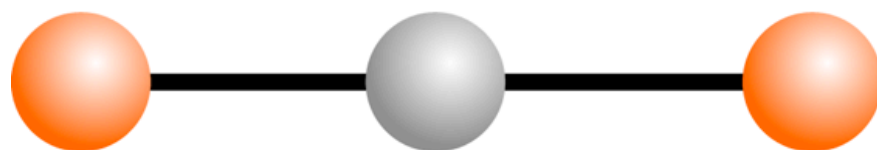
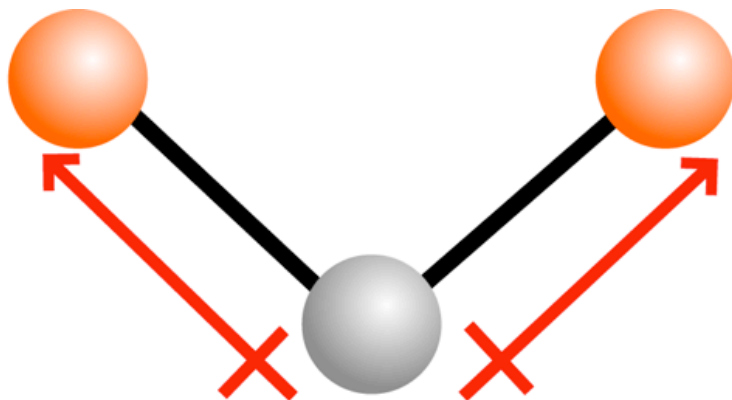


Optical selection rule 2 -- a change in molecular dipole moment ( $\Delta\mu/\Delta x$ ) must occur with the vibrational motion of a mode.

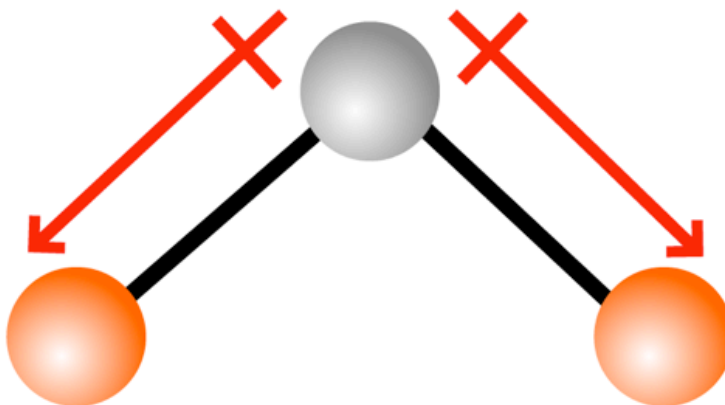
Let's break down the antisymmetric CO<sub>2</sub> stretch to see if the molecular dipole moment varies during this vibration.



...what about the bend?



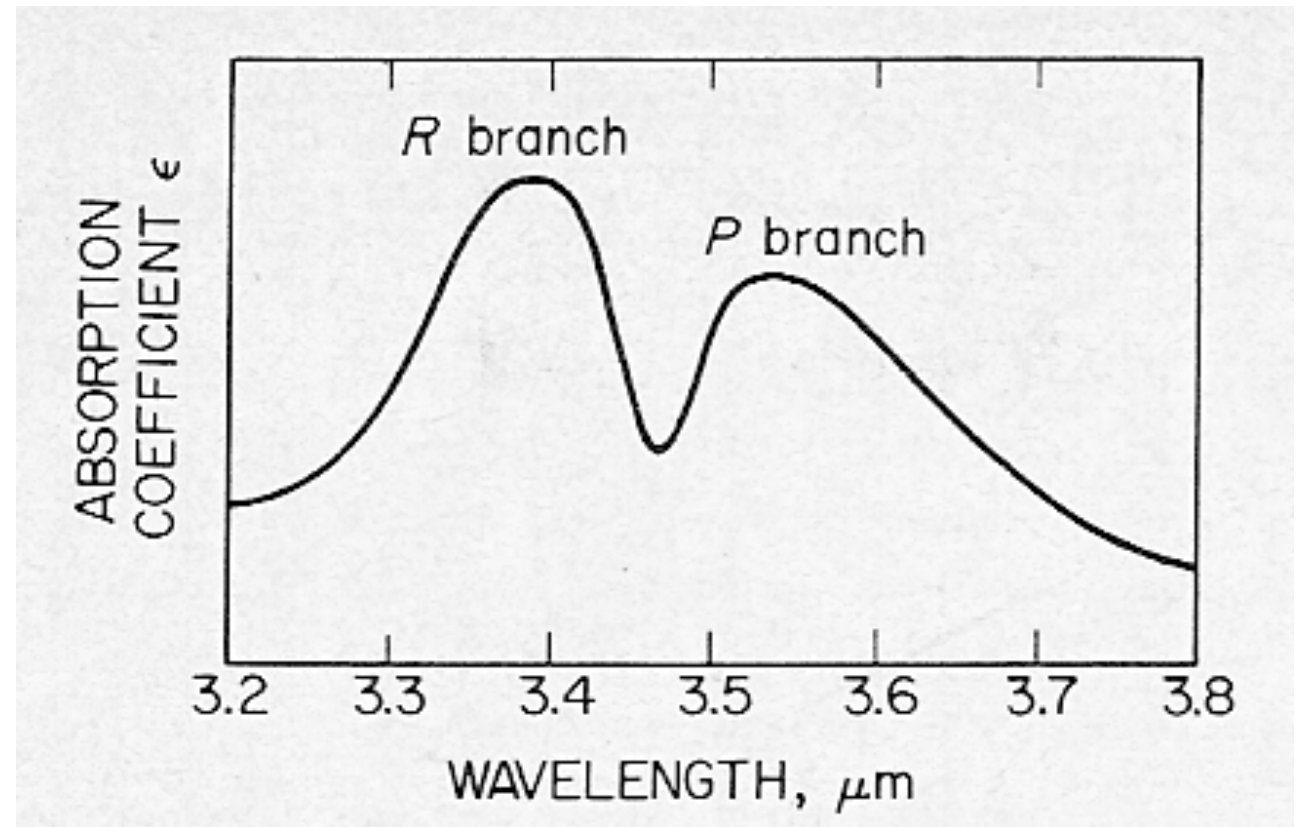
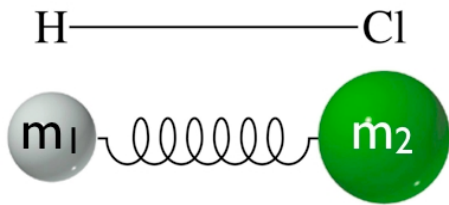
$= 0 \quad \nu = 667 \text{ cm}^{-1}$



## Summary:

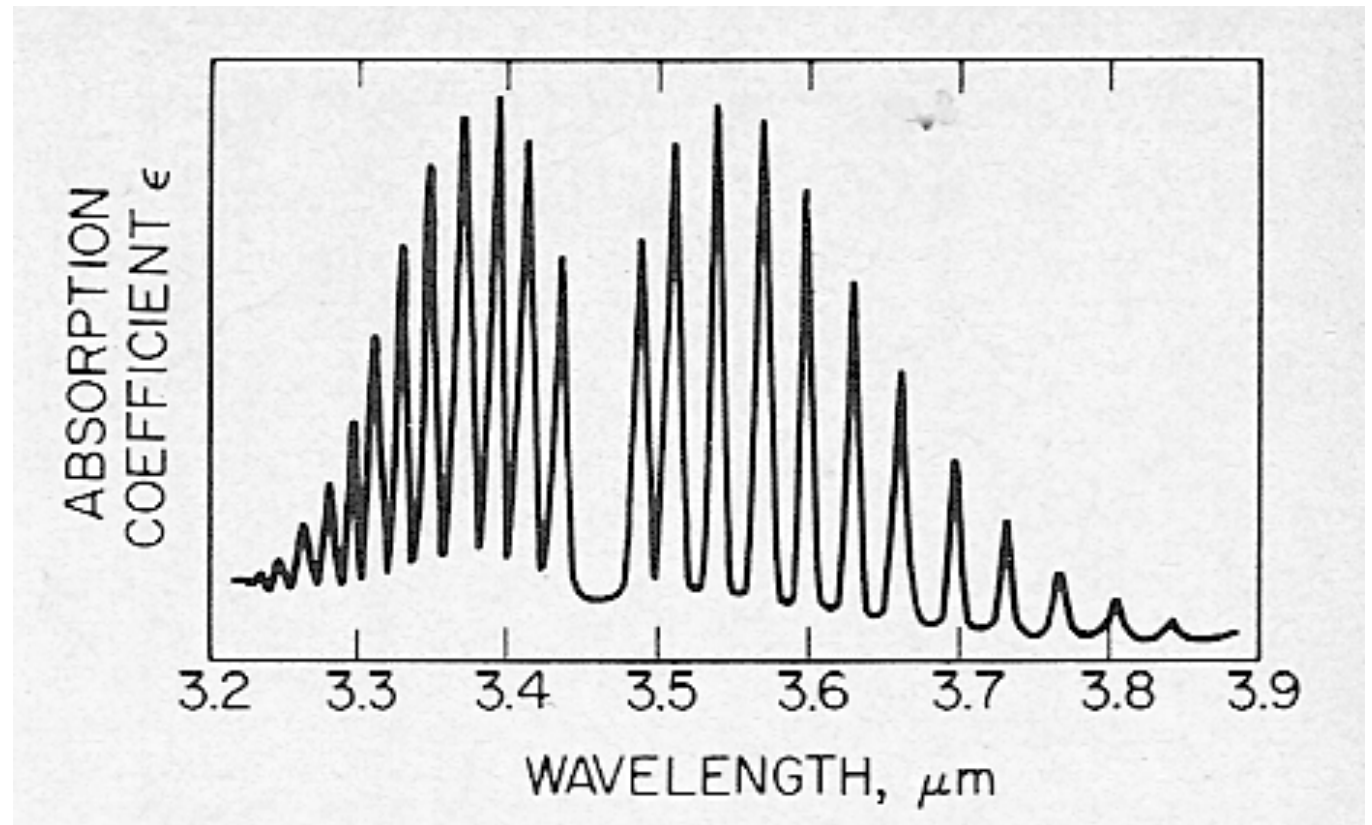
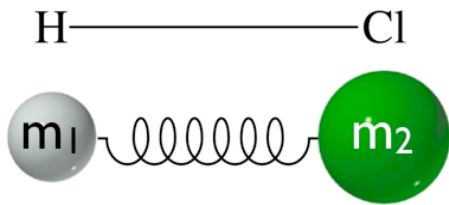
- ... CO<sub>2</sub> has  $3N-5 = 4$  vibrational modes:
  - 2 “degenerate” bends ( $\nu = 667 \text{ cm}^{-1}$ )
  - 1 asymmetric stretch. ( $\nu = 2349 \text{ cm}^{-1}$ )
  - 1 symmetric stretch. ( $\nu = 1380 \text{ cm}^{-1}$ )
- ...Three of these modes alter the dipole moment of the molecule (both bends and the asym stretch); one does not (the sym stretch).
- ...This means that CO<sub>2</sub> can absorb and radiate IR light at energies corresponding to  $667 \text{ cm}^{-1}$  and  $2349 \text{ cm}^{-1}$ ,  
**even though it has no permanent dipole moment.**

# HCl IR absorption spectrum in the gas phase:



In the gas phase, the lineshape of the IR absorption bands molecules becomes broad.

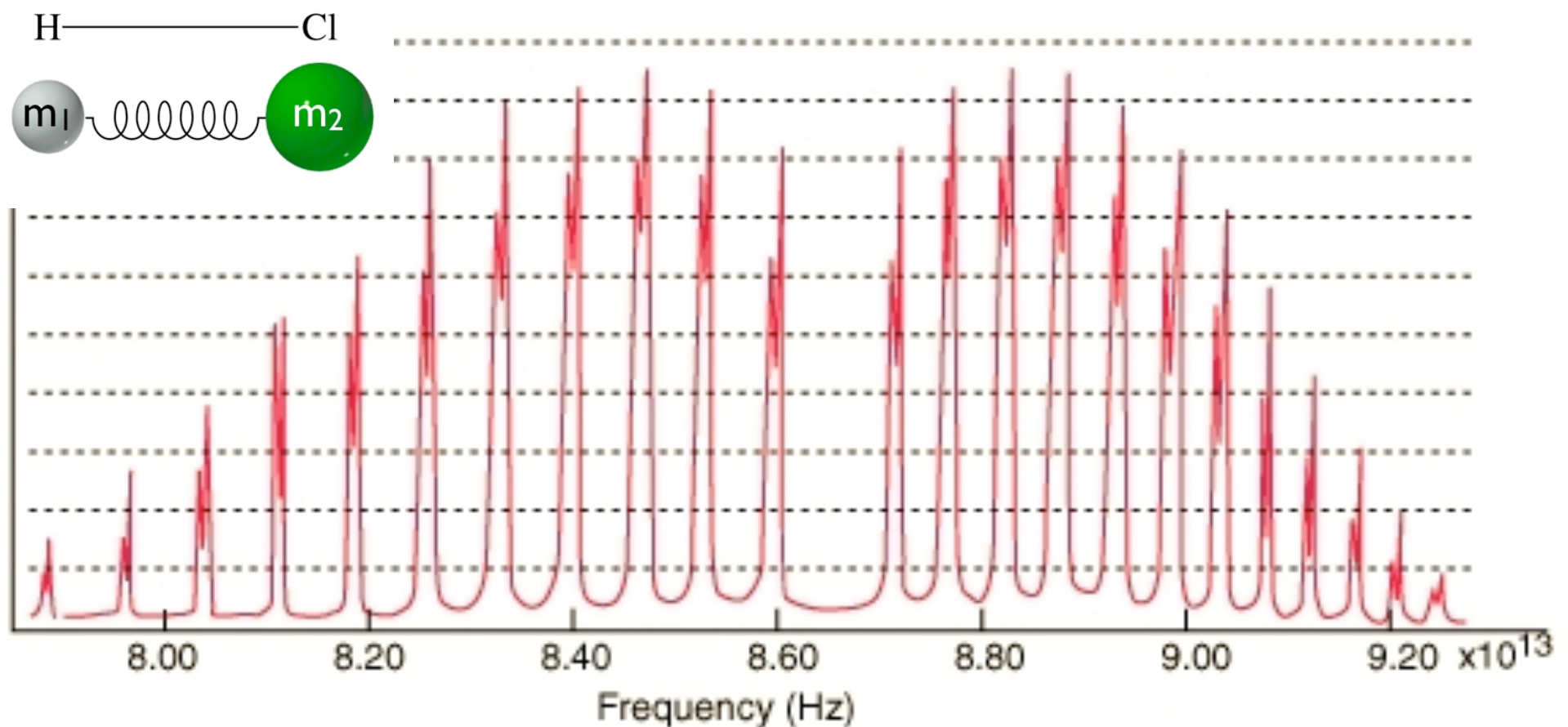
# HCl IR absorption spectrum in the gas phase:



At high resolution, MANY lines are observed. This splitting of the vibrational spectrum is due to the presence of many rotational quantum states.

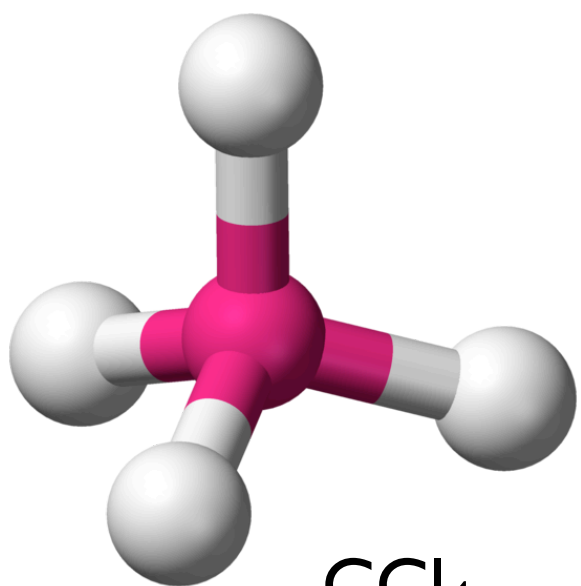


# HCl IR absorption spectrum in the gas phase:

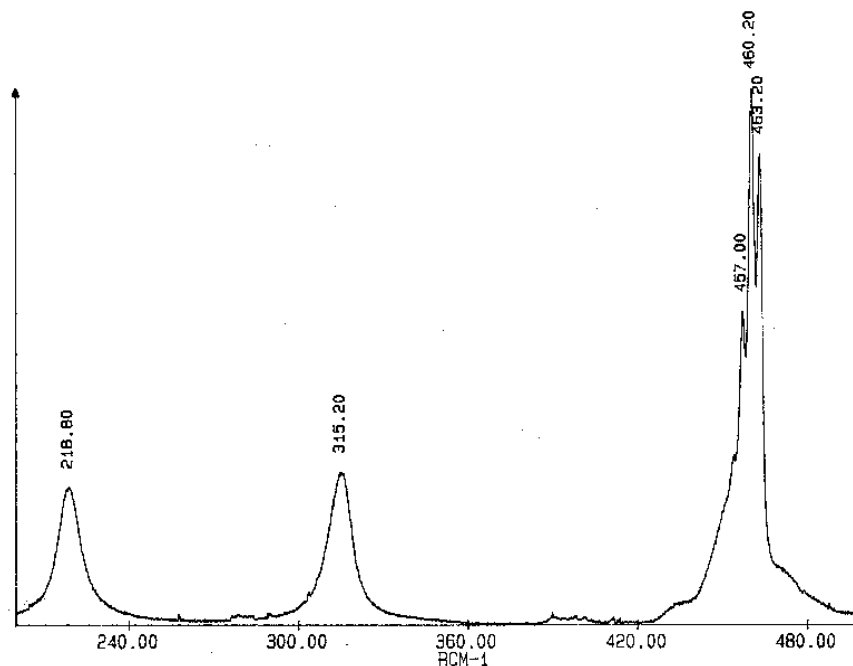


BTW -- a very high resolution spectrum of HCl shows PAIRS of bands due to  $\text{H}^{35}\text{Cl}$  and  $\text{H}^{37}\text{Cl}$ .

Symmetric stretches ARE NOT observed in the IR spectrum. But they ARE observed in another type of vibrational spectroscopy: **Raman scattering**

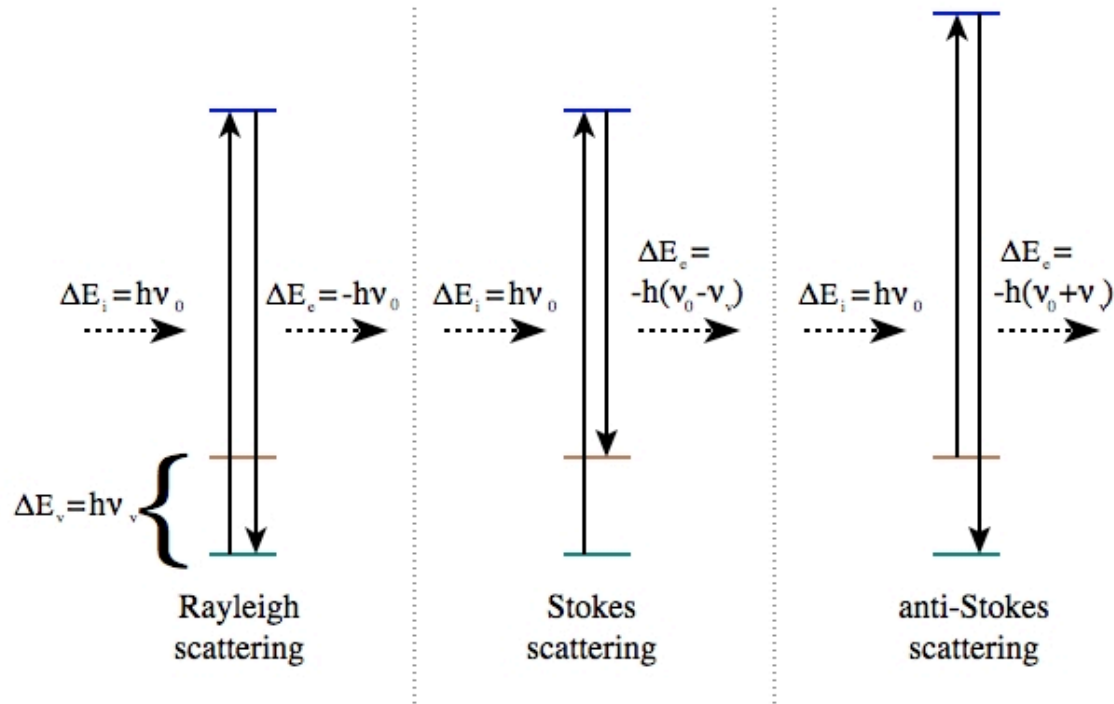


CCl4



CCl4 Raman Spectrum

# Raman scattering: Inelastic light scattering



C.V. Raman  
1888-1970

1930 Nobel Prize in Physics

# Raman scattering: Inelastic light scattering

Vibrational Raman scattering occurs when incident light **induces** a molecular dipole moment via the **molecular polarizability**. This dipole moment radiates light at a new wavelength, as determined by the energy levels of the molecule.



C.V. Raman  
1888-1970

1930 Nobel Prize in Physics

# Vibrational Raman scattering

Induced molecular dipole moment:

$$\mu = \alpha E$$

$\alpha$  is the molecular polarizability.

$$E(t) = E_0 \cos(\omega_0 t)$$

$$\alpha(t) = \alpha_0 + \alpha_{vib} \cos(\omega_{vib} t)$$

$$\mu(t) = \alpha(t) E(t)$$

$\omega_0$  is a visible frequency

$\omega_{vib}$  is a vibrational (IR) frequency

## Vibrational Raman scattering - Classical description

$$\mu(t) = \alpha_0 E_0 \cos(\omega_0 t) + \frac{1}{2} \alpha_{vib} E_0 \left( \left[ \cos(\omega_0 + \omega_{vib}) t \right] + \left[ \cos(\omega_0 - \omega_{vib}) t \right] \right)$$

Three colors are scattered from the molecules:

$\omega_0$  is called Rayleigh Scattering

$\omega_0 - \omega_{vib}$  is called Stokes Raman Scattering

$\omega_0 + \omega_{vib}$  is called Anti-Stokes Raman Scattering

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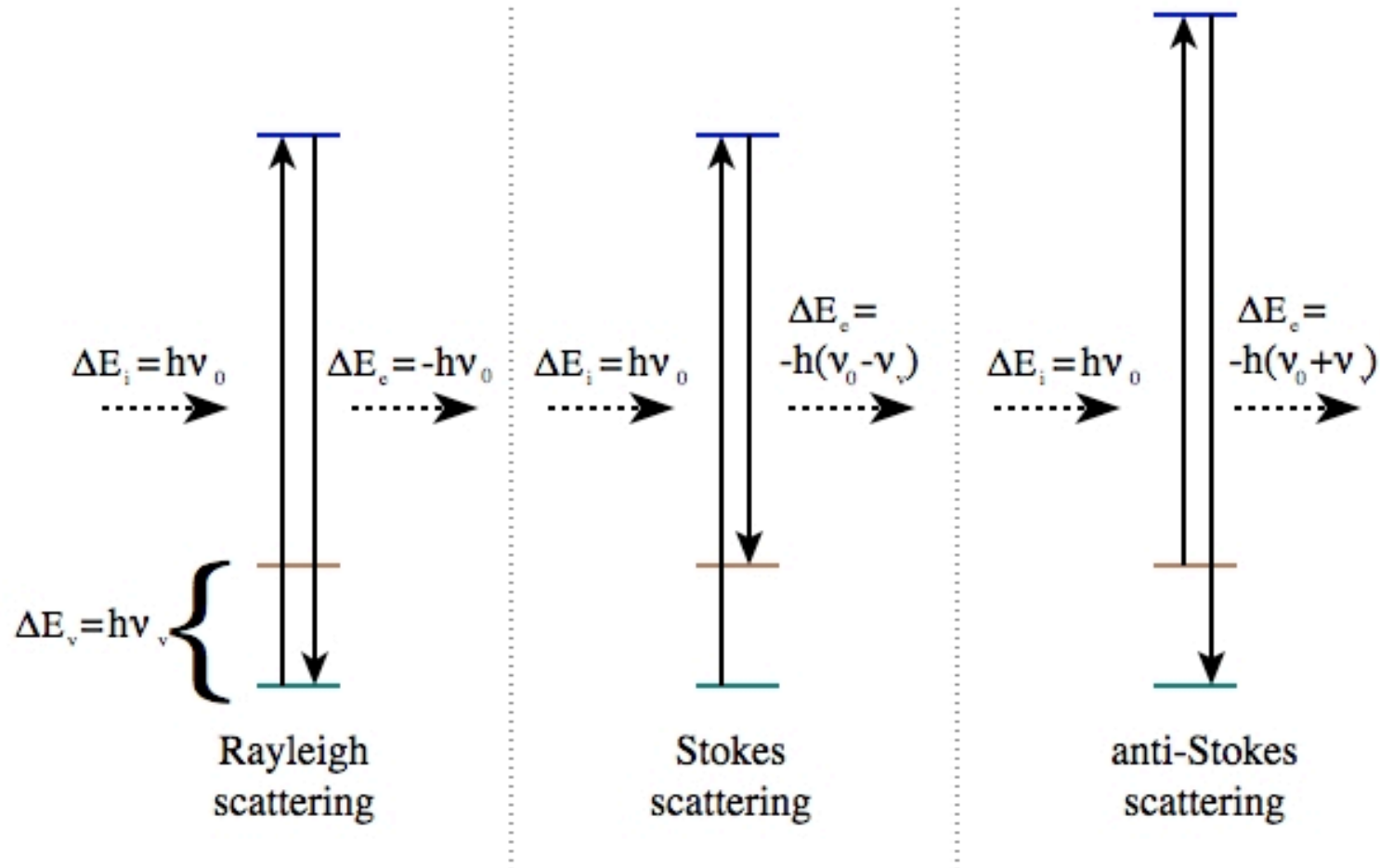
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# Vibrational Raman scattering - QM description



$\omega_0$  is called Rayleigh Scattering

$\omega_0 - \omega_{\text{vib}}$  is called Stokes Raman Scattering

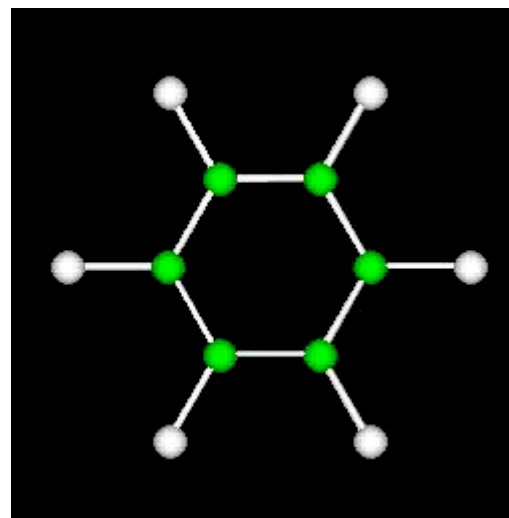
$\omega_0 + \omega_{\text{vib}}$  is called Anti-Stokes Raman Scattering

# Vibrational Raman scattering selection rules

$$\mu(t) = \alpha_0 E_0 \cos(\omega_0 t) + \frac{1}{2} \alpha_{vib} E_0 \left( \left[ \cos(\omega_0 + \omega_{vib}) t \right] + \left[ \cos(\omega_0 - \omega_{vib}) t \right] \right)$$

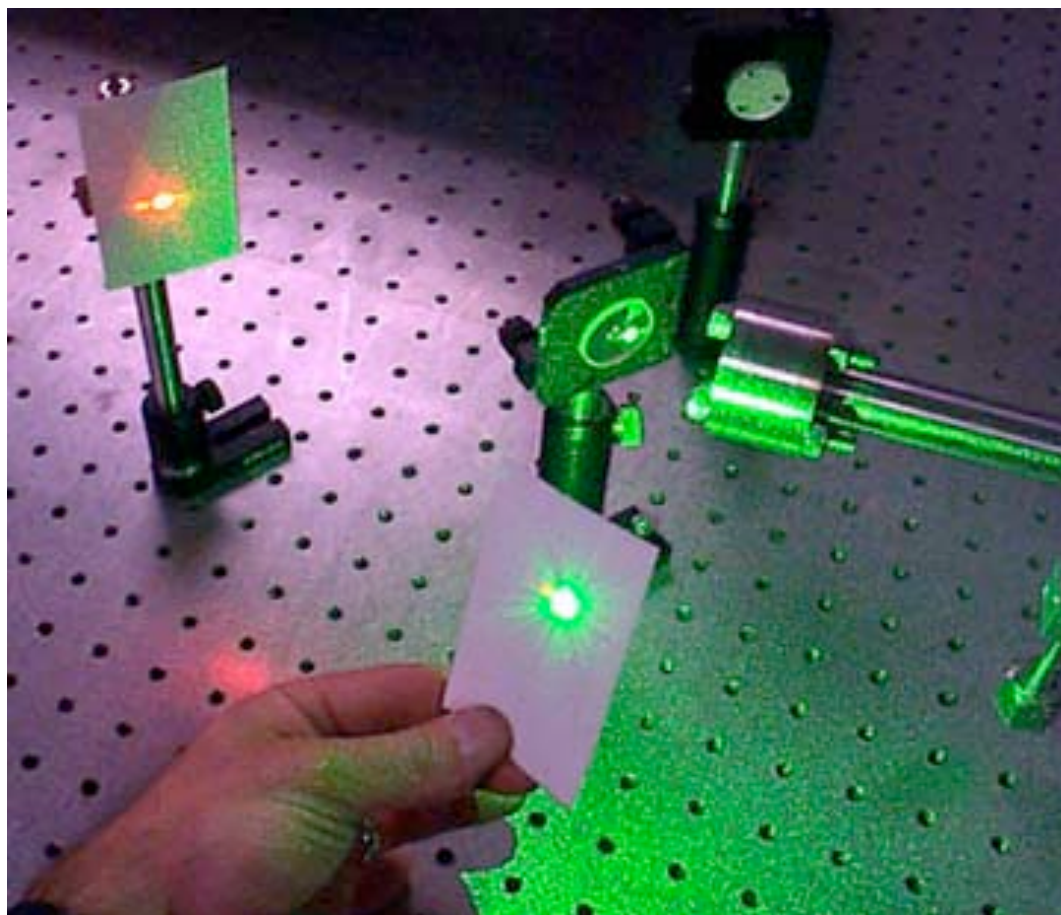
- For a normal mode vibration to be Raman active,  $\Delta\alpha/\Delta x$  must be non-zero: the vibration must change the molecular polarizability.

Symmetric modes can have large changes in molecular polarizability.



3000 cm<sup>-1</sup> CH breathing mode

# Vibrational Raman scattering



$$\lambda_0 = 500 \text{ nm}$$

$$\omega_0 = 20,000 \text{ cm}^{-1}$$

$$\omega_s = \omega_0 - \omega_{\text{vib}}$$

$$\omega_s = 20,000 - 3000 \text{ cm}^{-1}$$

$$\omega_s = 16,000 \text{ cm}^{-1}$$

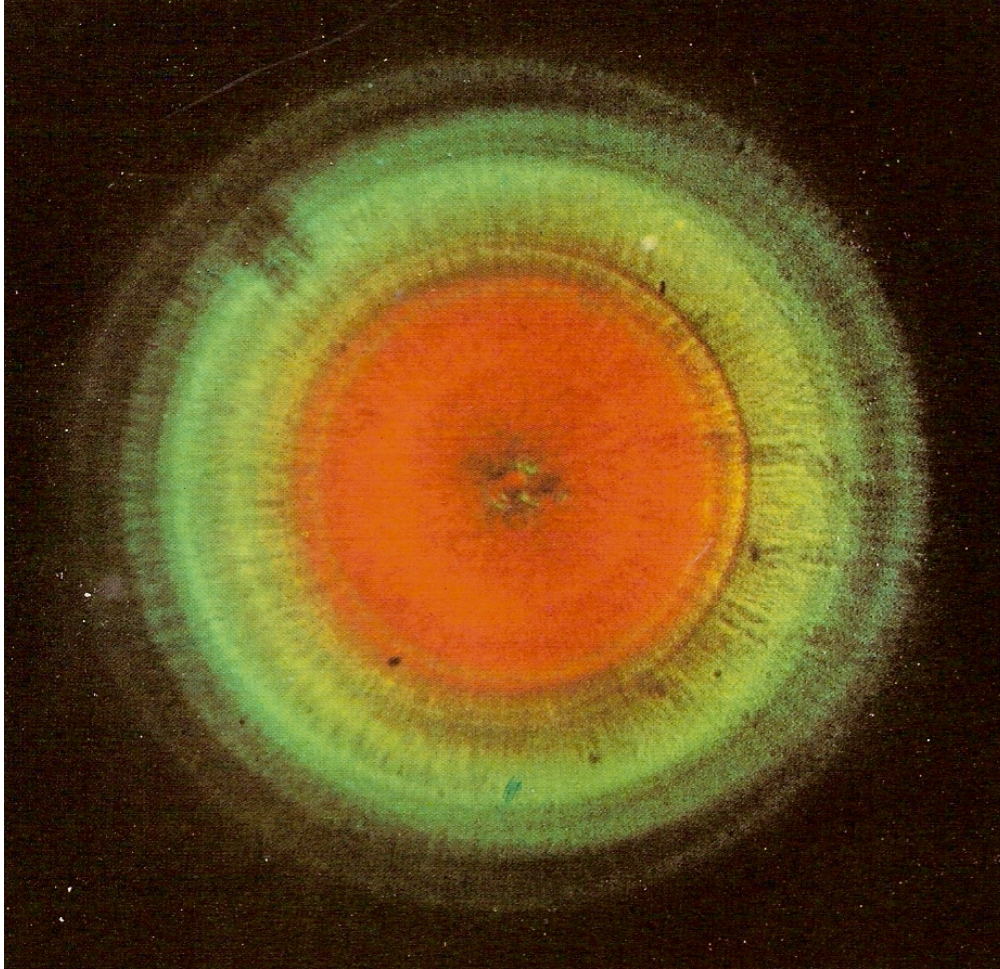
$$\lambda_s = 625 \text{ nm}$$

Benzene

$\omega_0 - \omega_{\text{vib}}$  is called Stokes Raman Scattering



# Vibrational Raman scattering



Benzene

$$\lambda_0 = 650 \text{ nm}$$

$$\omega_0 = 15,385 \text{ cm}^{-1}$$

$$\omega_{AS} = \omega_0 + \omega_{vib}$$

$$\omega_{AS} = 15,385 + 3000 \text{ cm}^{-1}$$

$$\omega_{AS} = 18,385 \text{ cm}^{-1}$$

$$\lambda_{AS} = 544 \text{ nm}$$

$\omega_0 + \omega_{vib}$  is called Anti-Stokes Raman Scattering